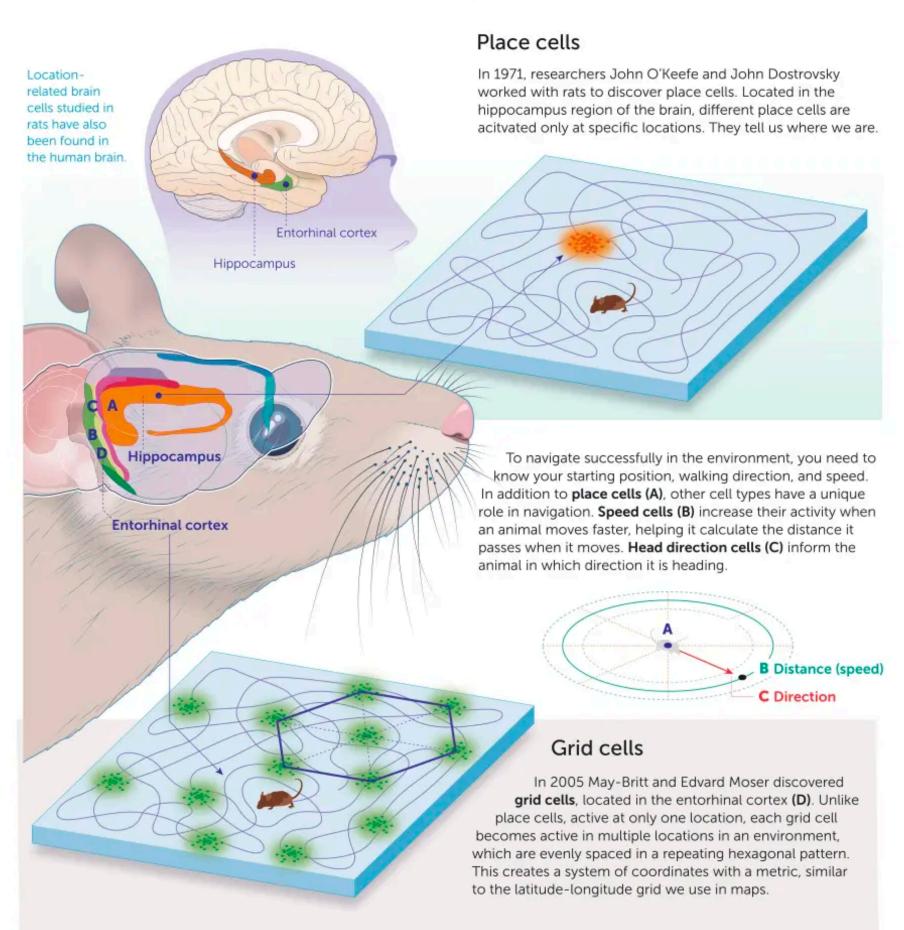
On Heads and Tails (2) Foundations seminar

Raul Vicente, 18th March 2025

🐉 frontiers

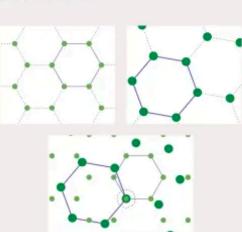
How our body's internal navigation system works

Different types of cells in the brain help us create an internal spatial representation of our environment or "cognitive map" and play a unique role in navigation.

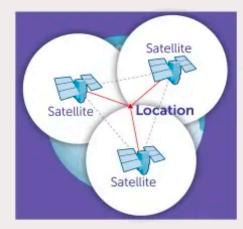


Grid cells help determine location ...

Different grid cells produce different hexagonal patterns at different scales and shifted with respect to the grids of other cells. Using the overlapping grids of several cells, location can be uniquely identified by coincident firing of multiple grid cells.



... similarly to GPS systems

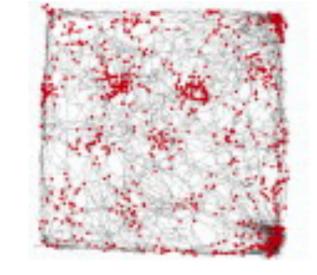


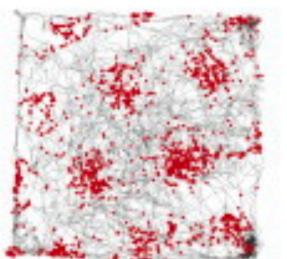
In a similar way, in GPS systems the distance measuring of one satelite alone can't pinpoint a location, but the overlap of at last three satelites is used to pinpoint the exact location of a receiver.

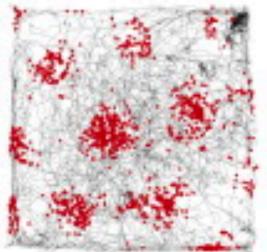
Source: How do we Find our Way? Grid Cells in the Brain, by May-Britt Moser. Published by Frontiers for Young Minds, Sep. 2021

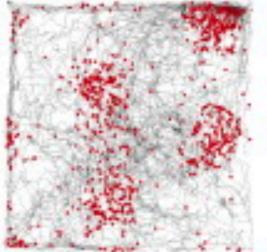
Infographic by 5W Infographic

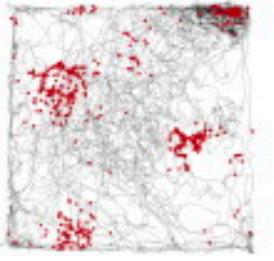




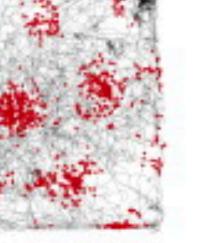


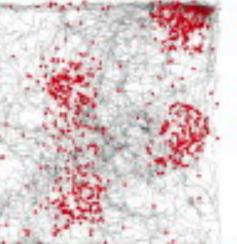


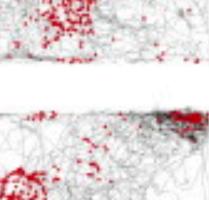


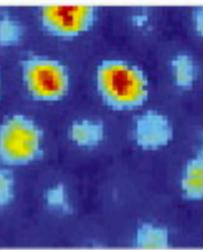


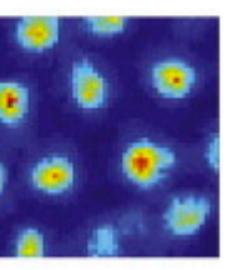


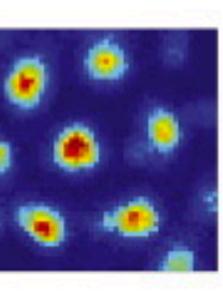


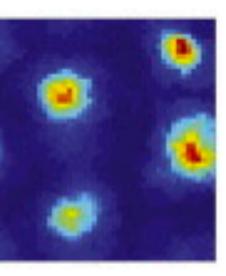


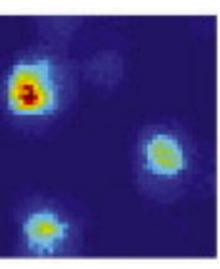


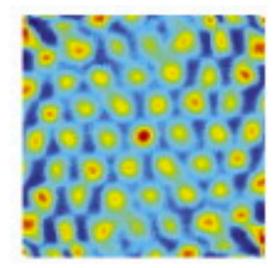


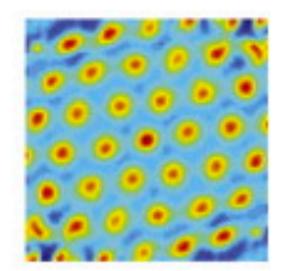


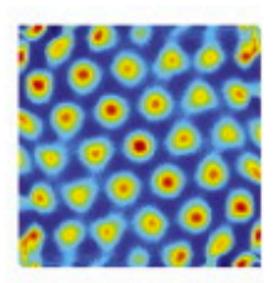


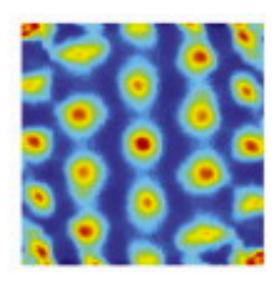


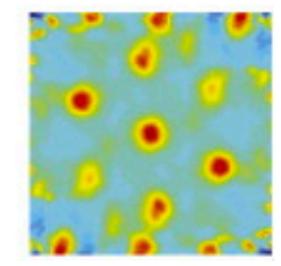


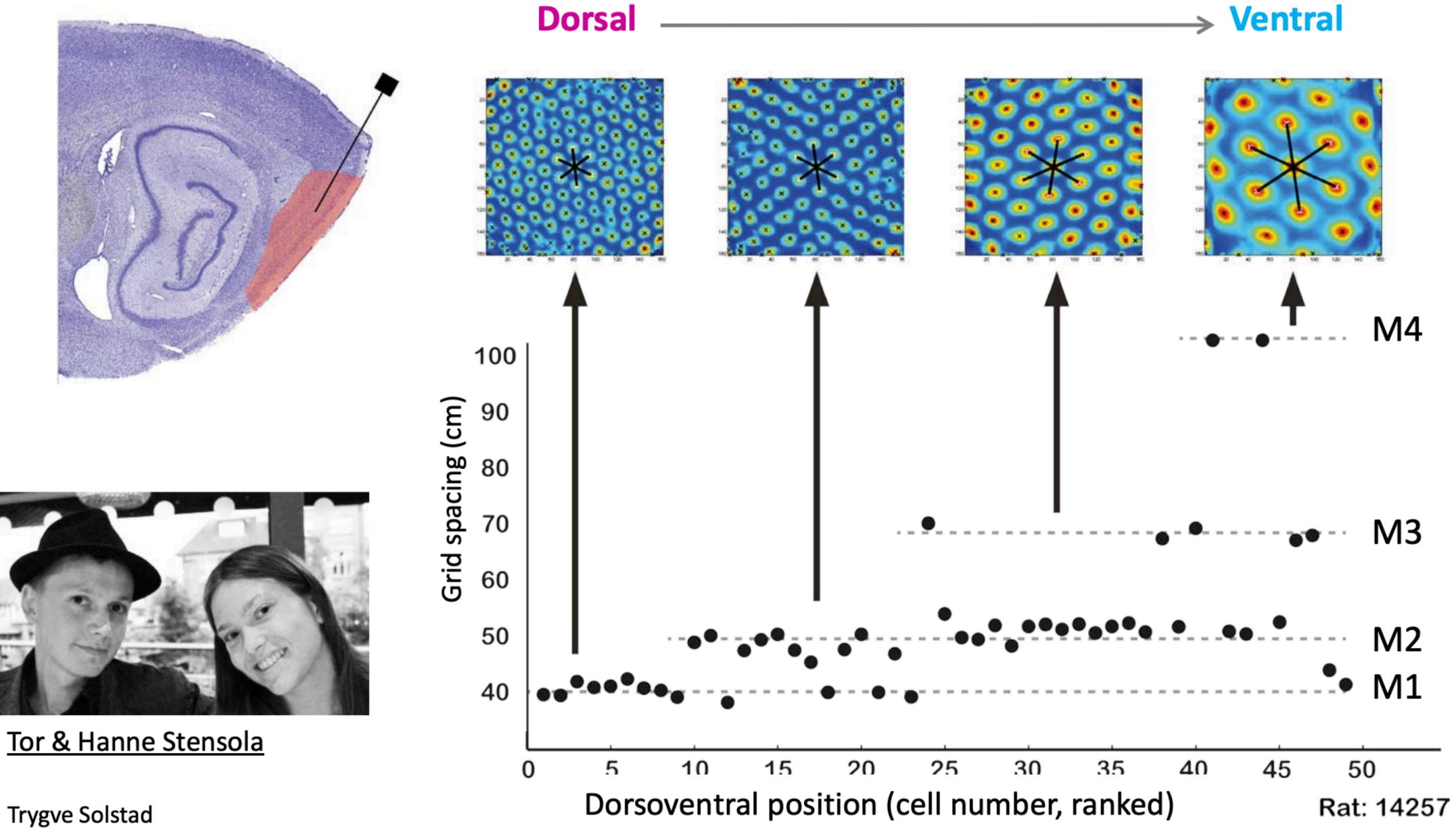












Trygve Solstad Kristian Frøland

Stensola et al. Nature, 492, 72-78 (2012)

Grid scale -> Eigenvalue associated to contributing eigenfunctions

Grid **pattern** -> Superposition of Fourier modes (**Eigenfunctions** of Laplacian **on the domain**)



1 How do these neurons activate on a hexagonal pattern?

2 What is their functional role?

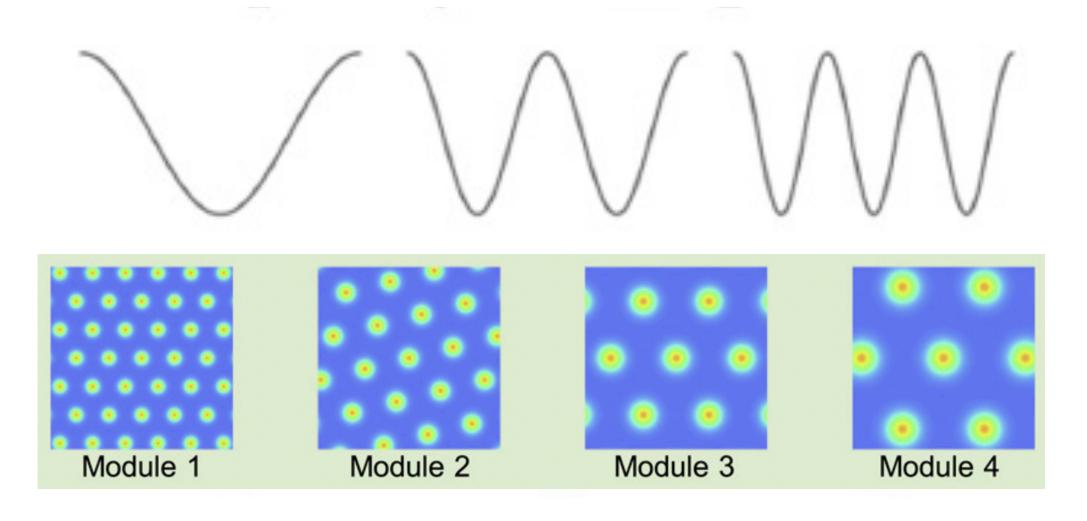


1 How do these neurons activate on a hexagonal pattern?

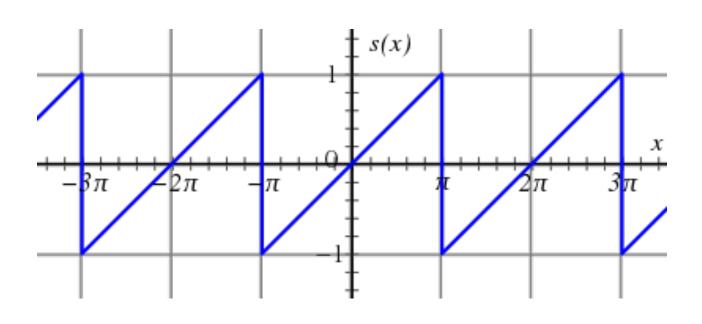
2 What is their functional role?

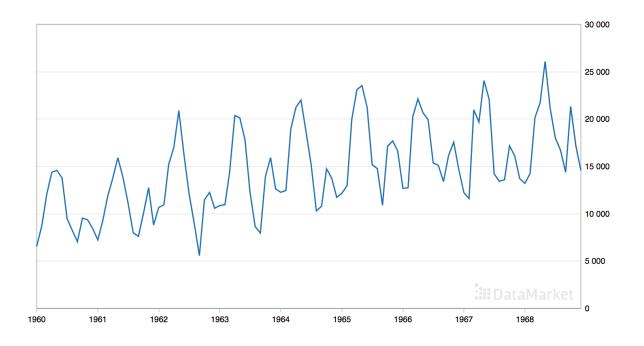
Research hypothesis

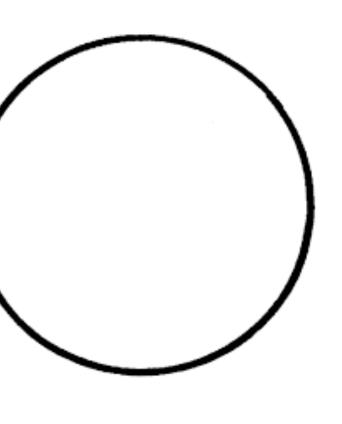
Grid cells = Fourier basis for the 2D environments



Fourier generalization to manifolds







$$\frac{d^2}{dx^2}\sin(wx) = -w^2\sin(wx)$$

$$\frac{d^2}{dx^2}\cos(wx) = -w^2\cos(wx)$$

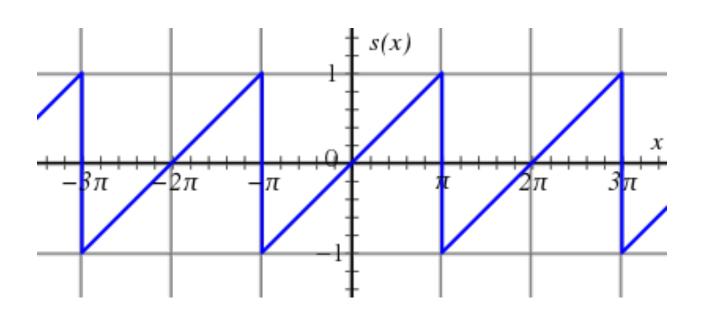
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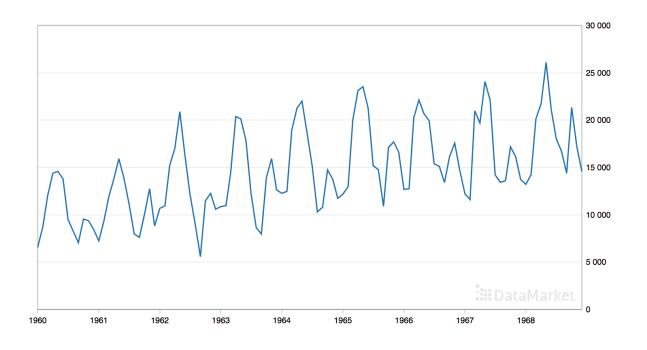
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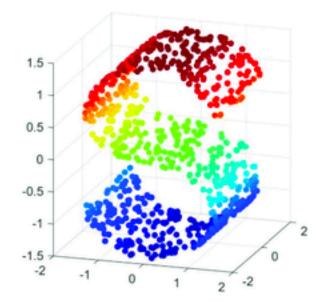


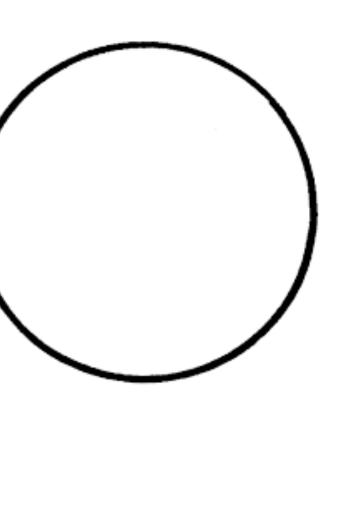


Fourier generalization to manifolds









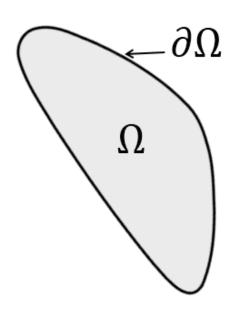
$$\frac{d^2}{dx^2}\sin(wx) = -w^2\sin(wx)$$

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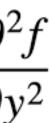
$$\frac{d^2}{dx^2}\sin(wx) = -w^2\sin(wx)$$

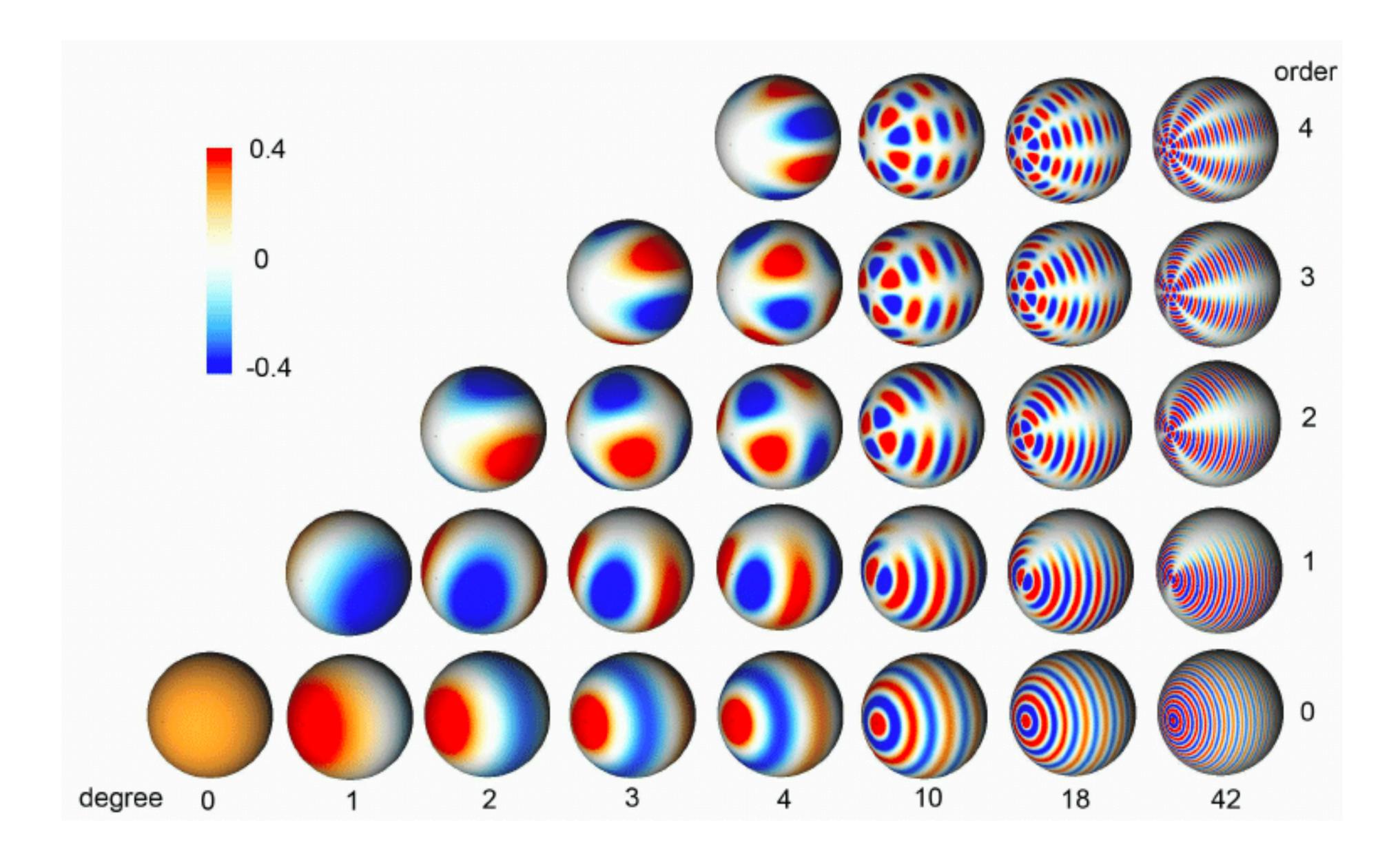
$$\frac{d^2}{dx^2}\cos(wx) = -w^2\cos(wx)$$

$$\nabla \cdot \nabla f = \nabla^2 f = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y}$$



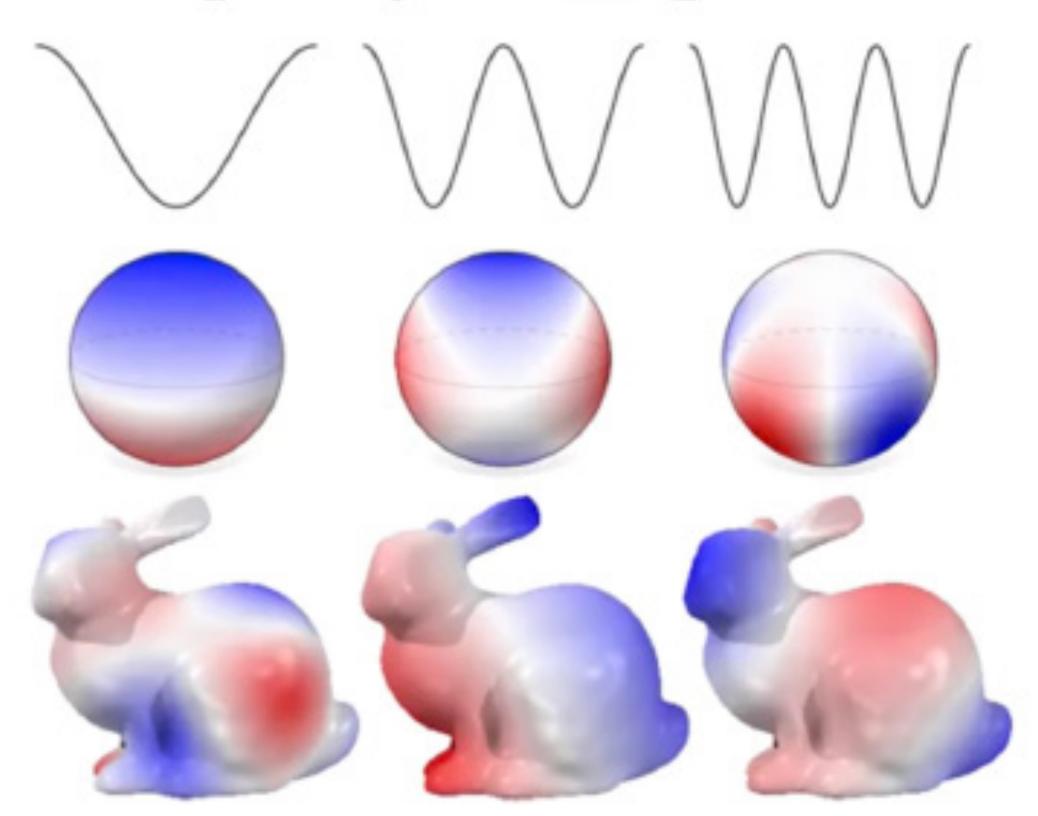






Generalised Fourier

frequency decomposition



 $\Delta \phi = \lambda \phi$

Basis set to represent functions

- Filtering of scales
- Smoothing and interpolation
- Spectral coordinates
- •

Representation Theory

Non-Abelian groups

Elements of the group are represented by invertible matrices and the group operation by matrix multiplication

Pontryagin Duality

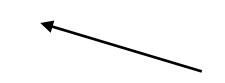
Functions on locally compact Abelian groups (\mathbb{Z}_2^n)

Extension of Fourier for functions on groups based on Pontryagin duality (group characters generalize complex exponentials as basis functions)

 $\hat{G} := \operatorname{Hom}(G, \mathbb{R}/\mathbb{Z})$

$$\chi: G \to U(1) \mid \chi(g_1g_2) = \chi(g_1)\chi(g_2)$$

$$\hat{f}(\chi) = \int_{G} f(x)\bar{\chi}(x)d\mu(x)$$



Spectral Geometry

Relation between geometry/topology and

Graphs

Relations between graph properties and spectra of the graph Laplacian

Hypercube

Fourier decomposition of Boolean functions: eigenvectors of Graph Laplacian are the characters of \mathbb{Z}_{2}^{n}

$$f: \{-1,1\}^n \to \mathbb{R}$$
$$f(x) = \sum_{S \subseteq [n]} \hat{f}(S) \chi_S(x), \quad \chi_S(x)$$

$$\hat{f}(S) = \left\langle f, \chi_S \right\rangle$$

Classical Fourier Analysis

Functions on $\mathbb{R} \ \mathbb{R}/\mathbb{Z} \ \mathbb{Z} \ \mathbb{Z}_n$

Representing functions as sum of trigonometric functions

$$\hat{f}(\xi) = \int_{-\infty}^{\infty} f(x)e^{-i2\pi\xi x}dx$$

eigenvalues and eigenfunctions of the Laplacian operator (Quantum mechanics)

Manifolds

 $(M,g) \rightarrow Spec(M,g)$

Eigenfunctions as minimizers of Dirichlet form

Flat torus (Lattices) \mathbb{R}^n / Δ

Poisson type formula relating the norms of lattice points (lengths of closed geodesics) to frequencies (eigenvalues)

Geodesic Flows

Relation between geometry/topology a and periodic geodesics (Classical mechanics)

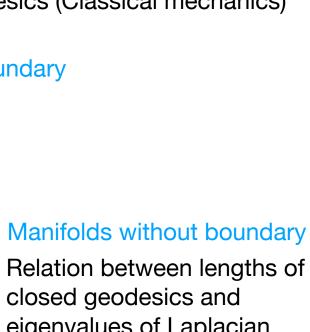
Manifolds with boundary

Billiard dynamics

closed geodesics and eigenvalues of Laplacian

 $Y_S(x) = \prod x_i$

$$\frac{1}{(4\pi t)^{d/2}} \operatorname{Vol}(\Delta) \sum_{\lambda \in \Delta} e^{-|\lambda|^2/4t} = \sum_{\iota \in \Delta^*} e^{-4\pi |\iota|^2 t}$$



Spectral geometry

Relation between geometry/topology and eigenvalues and eigenfunctions of the Laplacian operator

9.3.3 Facts

The Laplacian on any compact Riemannian manifold provides us with all the tools of Fourier analysis on our Riemannian manifold. Let us call a function ϕ an *eigenfunction* with *eigenvalue* the number λ if

 Δ

The set of all eigenvalues of Δ is an infinite discrete subset of \mathbb{R}^+ called the spectrum of Δ

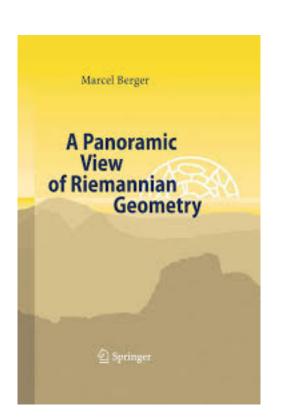
Spec $(M) = \{\lambda_k\}$

with λ_k tending to infinity with k. For each eigenvalue λ_i , the vector space of eigenfunctions ϕ satisfying

is always finite dimensional and its dimension is called the *multiplicity* of λ_i . Once we have a basis of the eigenfunctions with this eigenvalue written out, it is trivial to find an orthonormal basis

(where k runs from 1 to the multiplicity) of eigenfunctions. Here the orthonormalcy is to be understood for the global scalar product

$$\langle f,g \rangle_{L^2(M)} = \int_M fg \; .$$



$$f = \lambda f$$
.

$$= \{0 < \lambda_1 < \lambda_2 < \ldots\}$$

$$(9.5)$$

 $\Delta f = \lambda_i f$

 $\{\phi_k\}$

As for classical Fourier series, any reasonable function

has Fourier coefficients

and f is recovered from these coefficients by the converging series

In the same spirit, the scalar product of two functions is the sum of products of their coefficients:

where

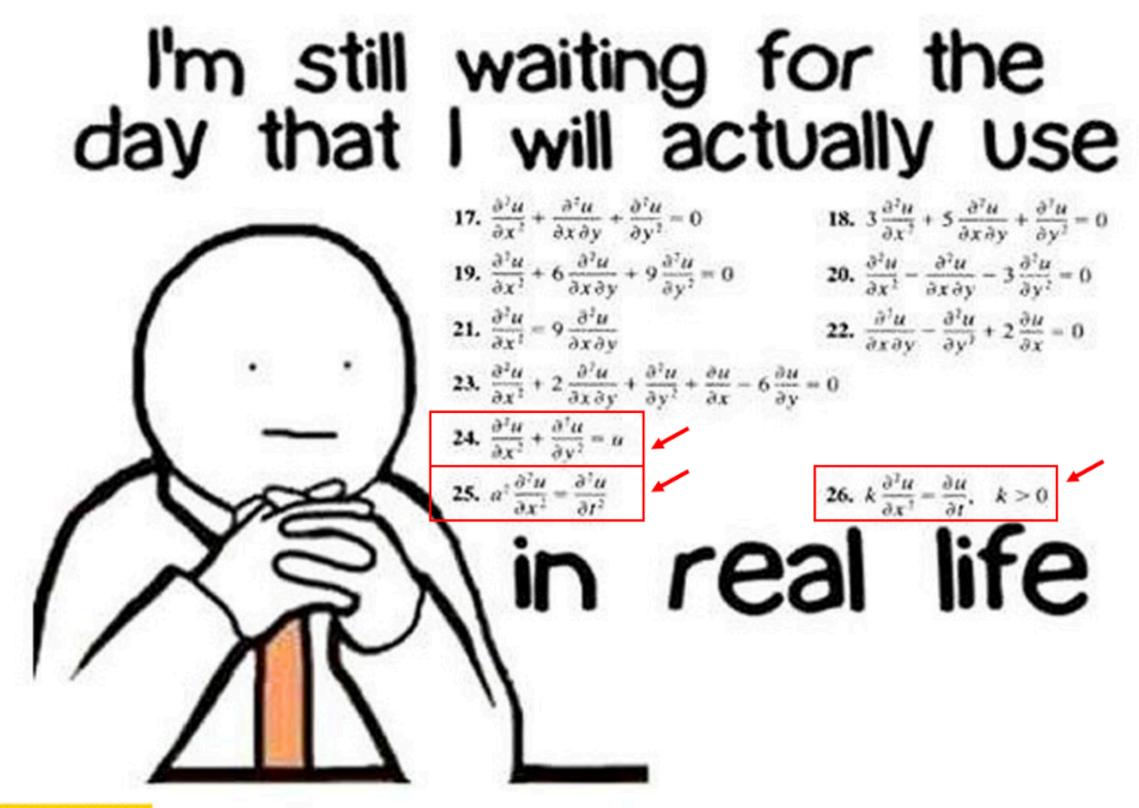
$$f: M \to \mathbb{R}$$

$$a_i = \int_M f\phi_i$$

$$f = \sum_i a_i \phi_i \; .$$

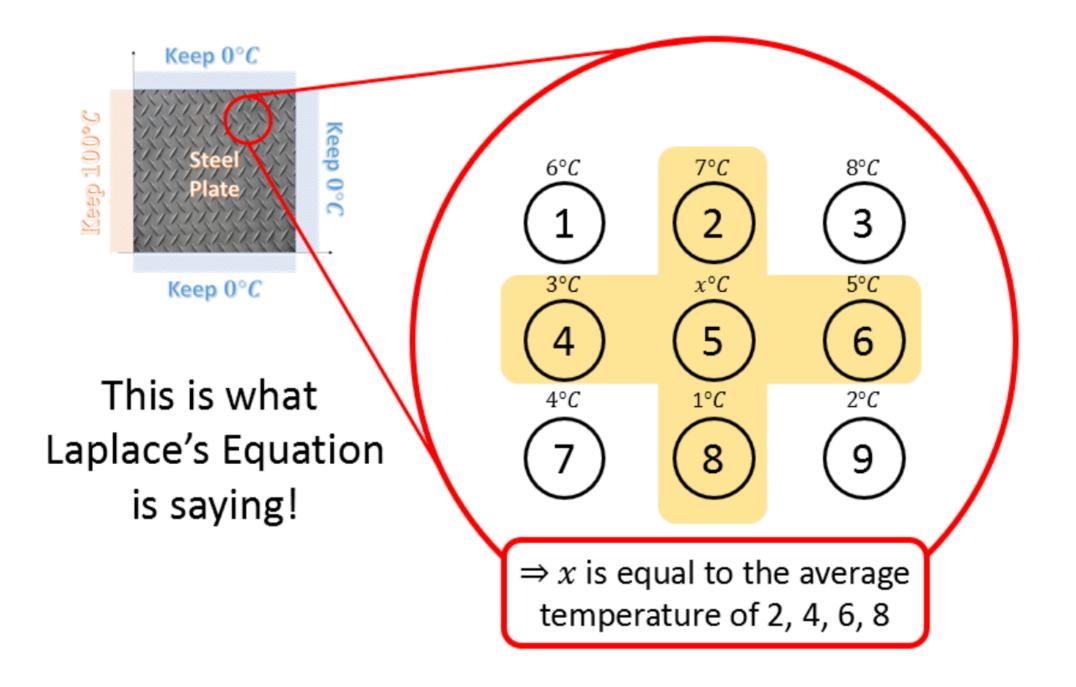
$$\int_{M} fg = \sum_{i} a_{i}b_{i}$$

$$f = \sum_{i} a_{i}\phi_{i}$$
$$g = \sum_{i} b_{i}\phi_{i} .$$



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$$\nabla \cdot \nabla f = \nabla^2 f = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2}$$



9.3.4 Heat, Wave and Schrödinger Equations

We will follow the same steps that we did in $\S1.8$: defining heat, wave and Schrödinger equations on Riemannian manifolds. The heat equation for the heat f(m,t) at time t at a point m of the Riemannian manifold M is

 Δf

point m is

The wave equation for the height f(m,t) of the "water" after time t at a $\Delta f = -\frac{\partial^2 f}{\partial t^2} \, .$ (9.7)

tion uses complex valued functions and is written

 $\hbar^2 \Delta$

where $i = \sqrt{-1}$ and \hbar is Planck's constant.

$$f = -\frac{\partial f}{\partial t} \ . \tag{9.6}$$

where if M were a surface, you would consider M covered in a thin sheet of water, or for M of three dimensions, M is a place through which sound is propagating. The wave equation can also be considered as describing the manifold M as a vibrating membrane object. Finally the Schrödinger equa-

$$\Delta f = i\hbar \frac{\partial f}{\partial t} \tag{9.8}$$

To solve these equations, at least formally, one uses the same trick as in §§1.8.1. To solve such an equation depending both on time t and a point $m \in$ M, the initial idea is to use the fact that, roughly by the Stone–Weierstraß approximation theorem, we need only to consider product functions

the heat equation precisely when the functions g and h satisfy

where

h

is the usual derivative.

Since the first fraction depends only on the point $m \in M$ and the second only on the time t their common value has to be a constant, call it λ . Then the function

$$h(t) = \begin{cases} e^{-\lambda t} & \text{fo} \\ e^{i\lambda t} & \text{fo} \\ e^{i\sqrt{\lambda}t} & \text{fo} \end{cases}$$

f(m,t) = g(m)h(t) .

One will subsequently consider series of them (as in the theory of Fourier series). Look for example at the heat equation. The function f = gh satisfies

$$\frac{\Delta g}{g} = -\frac{h'}{h} \tag{9.9}$$

$$a'(t) = \frac{dh}{dt}$$

 $g: M \to \mathbb{R}$

is an eigenfunction of Δ with eigenvalue λ , while h is an exponential decay at rate λ . If all eigenfunctions and eigenvalues of Δ are known, we can then solve the heat equation explicitly. Note that the time dependence h(t) is

> or the heat equation or the Schrödinger equation or the wave equation.

is to compute the Riemannian Fourier series

f(m,0)

and then

$$f(m,t) = \sum_{k=1}^{\infty} a_k \phi_k(m) e^{-\lambda_k t}$$

For the wave equation, the fundamental solution similar to equation 9.10 requires imaginary terms, i.e.

which are linear combinations of

$$\cos\left(\sqrt{\lambda_k}t\right)$$
 and $\sin\left(\sqrt{\lambda_k}t\right)$

But the dramatic difference between the heat equation and the wave equation is that waves demand not converging series, but distributions. Heat spreads out uniformly with time, while waves bounce up and down forever. This ma-

Another way to write the solution f(m, t) with initial temperature f(m, 0)

$$) = \sum_{k=1}^{\infty} a_k \phi_k$$

$$e^{i\sqrt{\lambda_k}t}$$

are

 $\sin \frac{\pi}{-}$

where m and n are any integers, yielding the set of eigenvalues

$$\lambda = \pi^2 \left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right).$$

 $m^2 + n^2 < \lambda^2$

is asymptotic to $\pi^2 \lambda^2 + O(\lambda^{\varepsilon+1/2})$ as $\lambda \to \infty$, for any $\varepsilon > 0$, but there is still no proof. This is called the Gauß circle problem; see §§1.8.5. However, it is known that $\pi^2 \lambda^2 + O(\lambda^{1/2})$ is too small.

There are very few examples where the spectrum or the eigenfunctions can be determined explicitly. Two old standards are rectangles and disks. In both cases, separation of variables disentangles the eigenfunctions. Using again the Stone–Weierstraß theorem, and because the boundary condition agrees with the separation, on a rectangle one need only look for product functions f(x,y) = g(x)h(y), and there will be no other eigenfunctions. If the rectangle has side lengths a and b respectively, then the eigenfunctions

$$\frac{mx}{a}\sin\frac{\pi ny}{b} \tag{1.21}$$

i.e. to obtain an eigenvalue at most λ we need m and n to be integer points inside a certain ellipse. We will see below that a simple expression yields an easy first order approximation of $N(\lambda)$ when $\lambda \to \infty$, but the second order term in λ is related to deep number theory and is still not completely understood today. It is believed that the number of integers m, n with

How about a second order approximation? In 1954, Pleijel got the next order approximation. In his paper Kac 1966 [775], Kac works quite hard to get the third term, guessing that the right formula should be:

$$\sum_{i=1}^{\infty} e^{-\lambda_i t} \sim \frac{\operatorname{Area}(D)}{2\pi t} - \frac{\operatorname{length}(\partial D)}{\sqrt{2\pi t}} + \frac{1}{6}(1-r)$$

where r is the number of holes inside D. The second term is Pleijel's. Note that can one hear the area and the perimeter of D, hence the isoperimetric inequality yields again the fact that disks are characterized by their spectrum. Kac could only prove the third term for polygons. It was proven the next year, 1967, in the very general context of Riemannian manifolds with boundary by McKean and Singer in their fundamental paper of 1967 [910]. You will read much more about it in chapter 9.

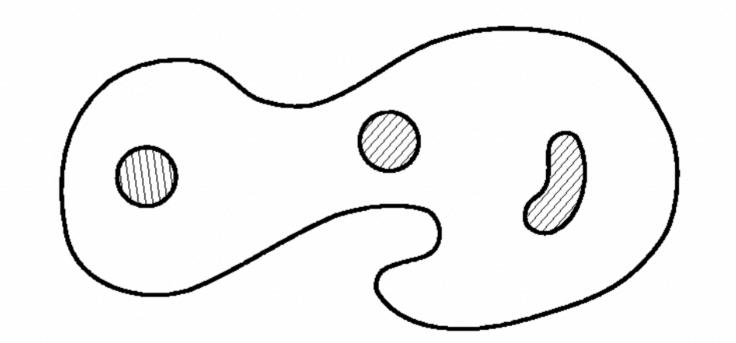


Fig. 1.97. One can hear the number of holes

The Tauberian theorem above shows that the first terms of $N(\lambda)$ and $\sum_{i=1}^{\infty} \exp(-\lambda_i t)$, can each be acquired from the other. But this does not



Can One Hear the Shape of a Drum? Author(s): Mark Kac Reviewed work(s): Source: The American Mathematical Monthly, Vol. 73, No. 4, Part 2: Papers in Analysis (Apr., 1966), pp. 1-23 Published by: Mathematical Association of America Stable URL: http://www.jstor.org/stable/2313748 Accessed: 15/03/2012 22:10

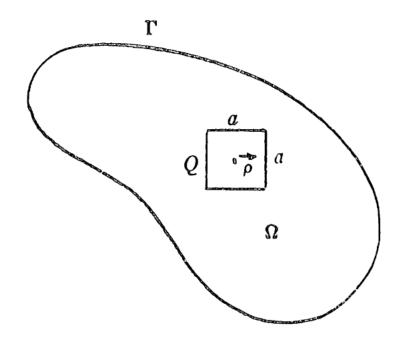
CAN ONE HEAR THE SHAPE OF A DRUM?

MARK KAC, The Rockefeller University, New York

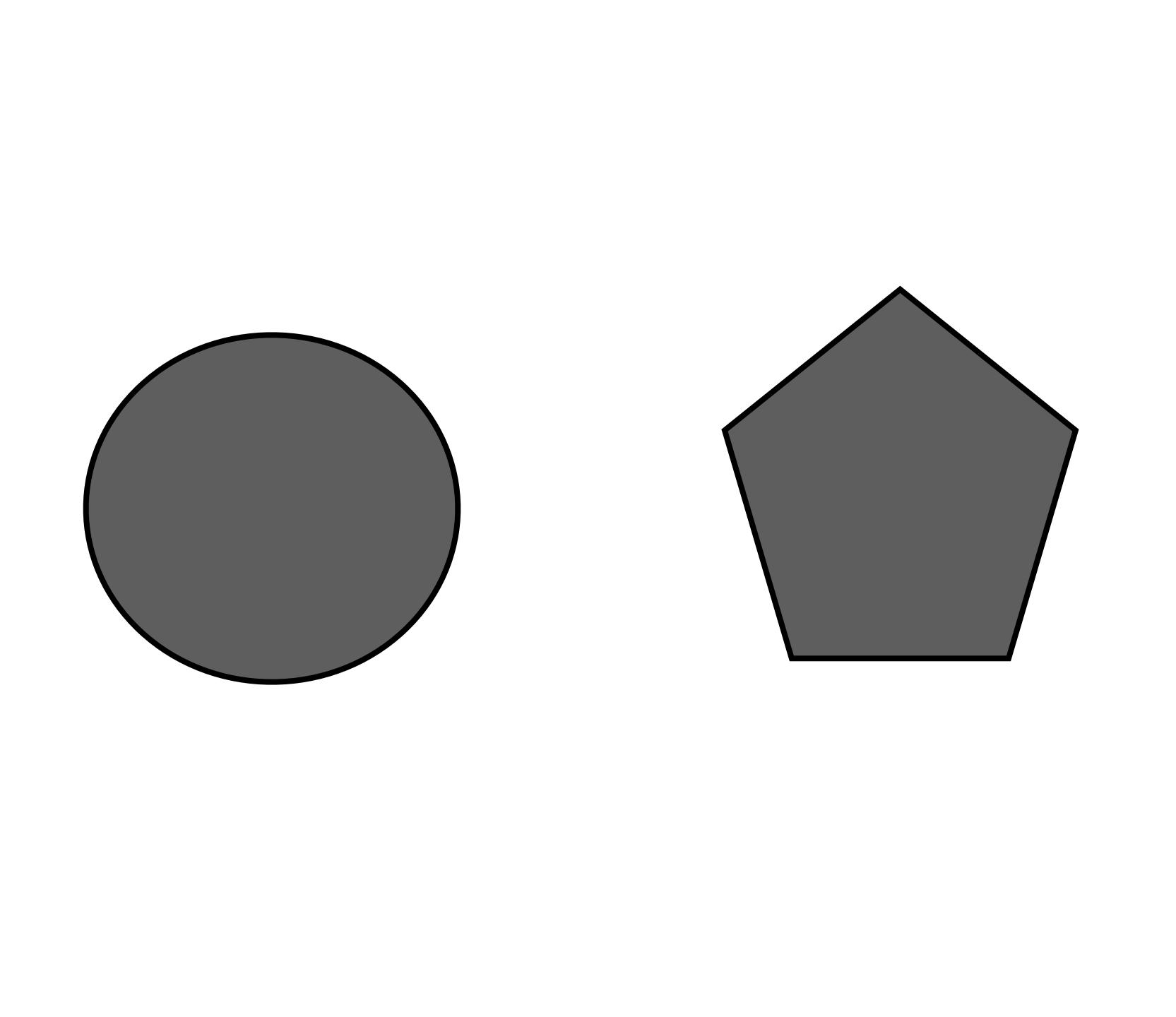
To George Eugene Uhlenbeck on the occasion of his sixty-fifth birthday

"La Physique ne nous donne pas seulement l'occasion de résoudre des problèmes . . . , elle nous fait presentir la solution." H. POINCARÉ.

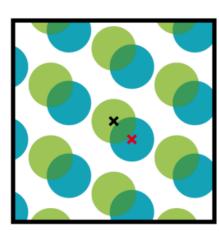
Before I explain the title and introduce the theme of the lecture I should like to state that my presentation will be more in the nature of a leisurely excursion than of an organized tour. It will not be my purpose to reach a specified destination at a scheduled time. Rather I should like to allow myself on many occasions the luxury of stopping and looking around. So much effort is being spent on streamlining mathematics and in rendering it more efficient, that a solitary transgression against the trend could perhaps be forgiven.



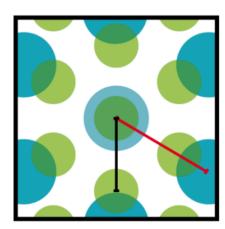




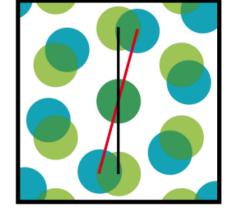
1. Phase

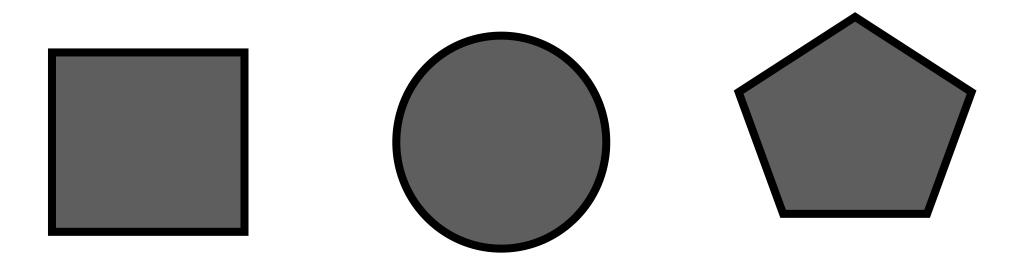


2. Scale









Faber-Krahn theorem

Rayleigh–Faber–Krahn inequality :=

Article View the content page [ctrl-option-c]

From Wikipedia, the free encyclopedia

In spectral geometry, the Rayleigh–Faber–Krahn inequality, named after its conjecturer, Lord Rayleigh, and two individuals who independently proved the conjecture, G. Faber and Edgar Krahn, is an inequality concerning the lowest Dirichlet eigenvalue of the Laplace operator on a bounded domain in \mathbb{R}^n , $n \geq 2$.^[1] It states that the first Dirichlet eigenvalue is no less than the corresponding Dirichlet eigenvalue of a Euclidean ball having the same volume. Furthermore, the inequality is rigid in the sense that if the first Dirichlet eigenvalue is equal to that of the corresponding ball, then the domain must actually be a ball. In the case of n=2, the inequality essentially states that among all drums of equal area, the circular drum (uniquely) has the lowest voice.



Grid scale -> Eigenvalue associated to contributing eigenfunctions

Grid **pattern** -> Superposition of Fourier modes (**Eigenfunctions** of Laplacian **on the domain**)



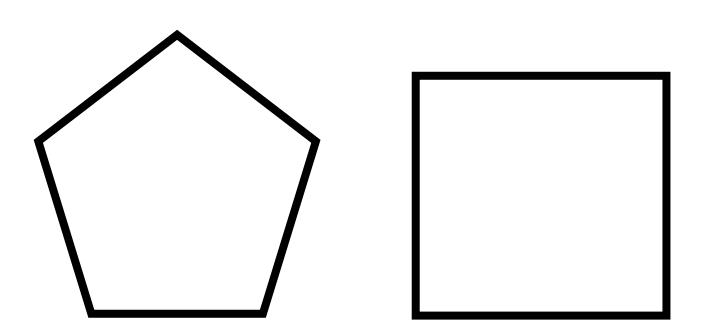
Rayleigh–Faber–Krahn inequality :=

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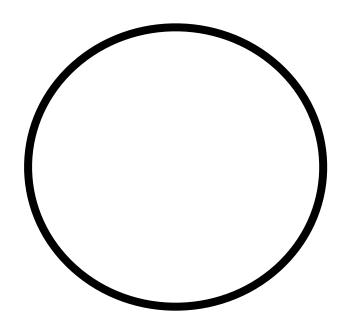
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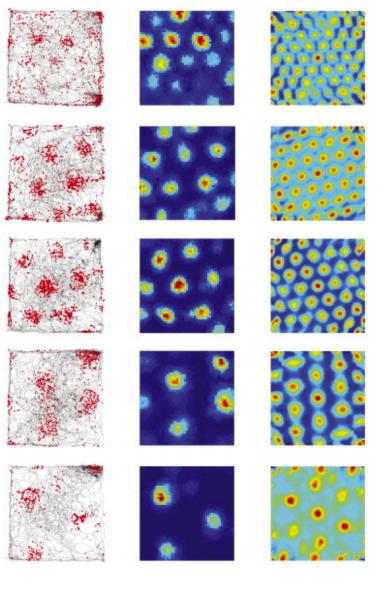
In spectral geometry, the Rayleigh–Faber–Krahn inequality, named after its conjecturer, Lord Rayleigh, and two individuals who independently proved the conjecture, G. Faber and Edgar Krahn, is an inequality concerning the lowest Dirichlet eigenvalue of the Laplace operator on a bounded domain in \mathbb{R}^n , $n \geq 2$.^[1] It states that the first Dirichlet eigenvalue is no less than the corresponding Dirichlet eigenvalue of a Euclidean ball having the same volume. Furthermore, the inequality is rigid in the sense that if the first Dirichlet eigenvalue is equal to that of the corresponding ball, then the domain must actually be a ball. In the case of n=2, the inequality essentially states that among all drums of equal area, the circular drum (uniquely) has the lowest voice.





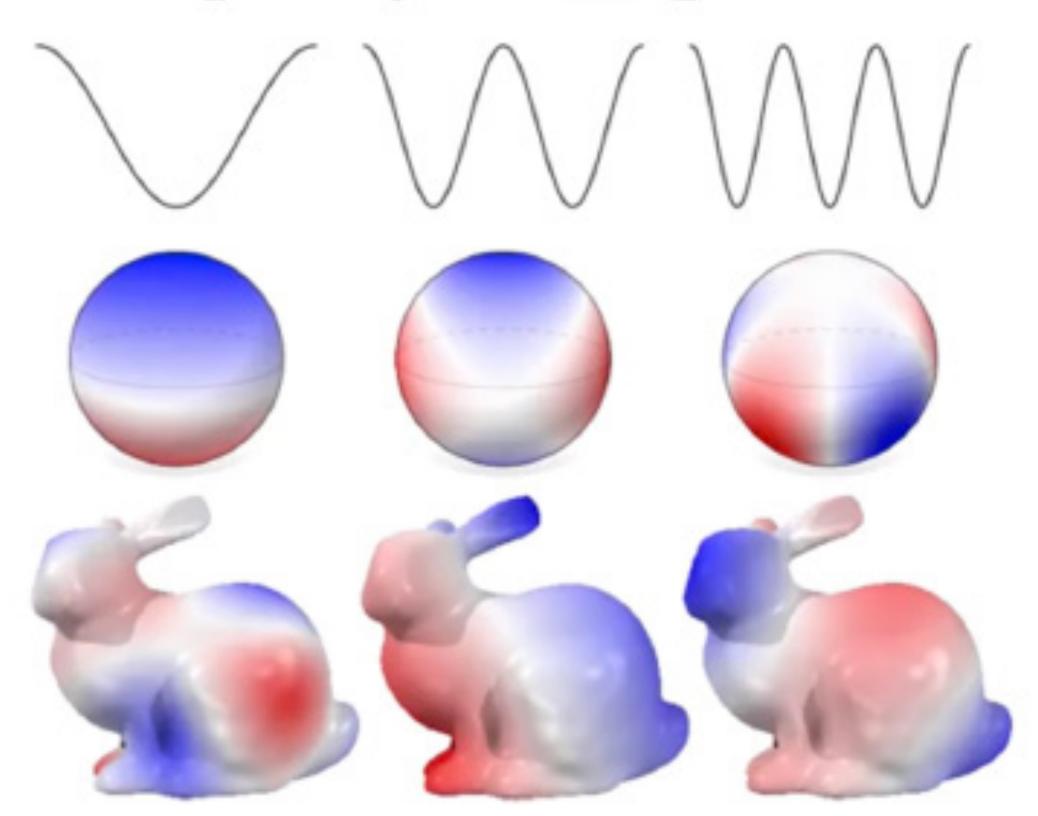






Generalised Fourier

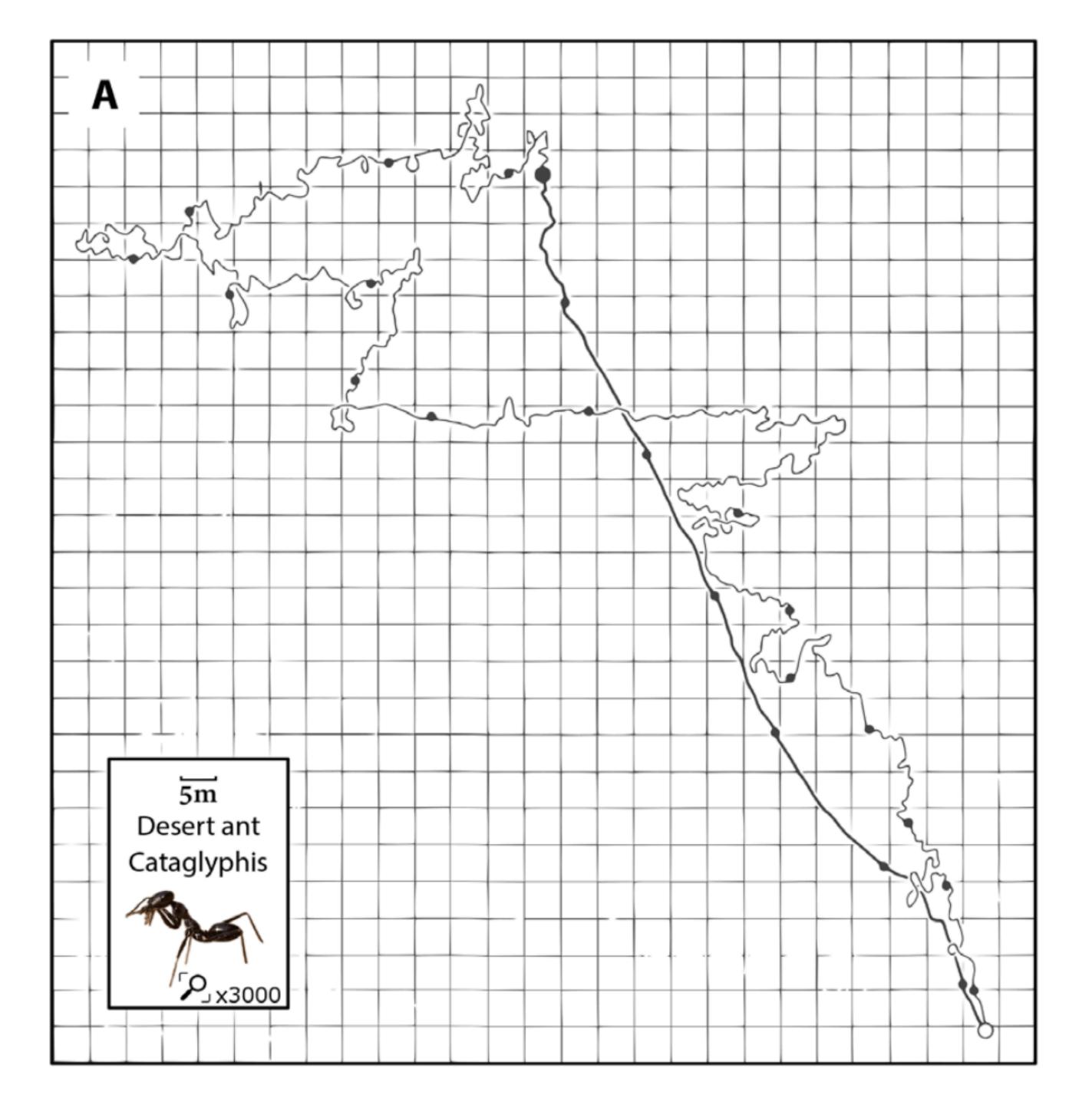
frequency decomposition



 $\Delta \phi = \lambda \phi$

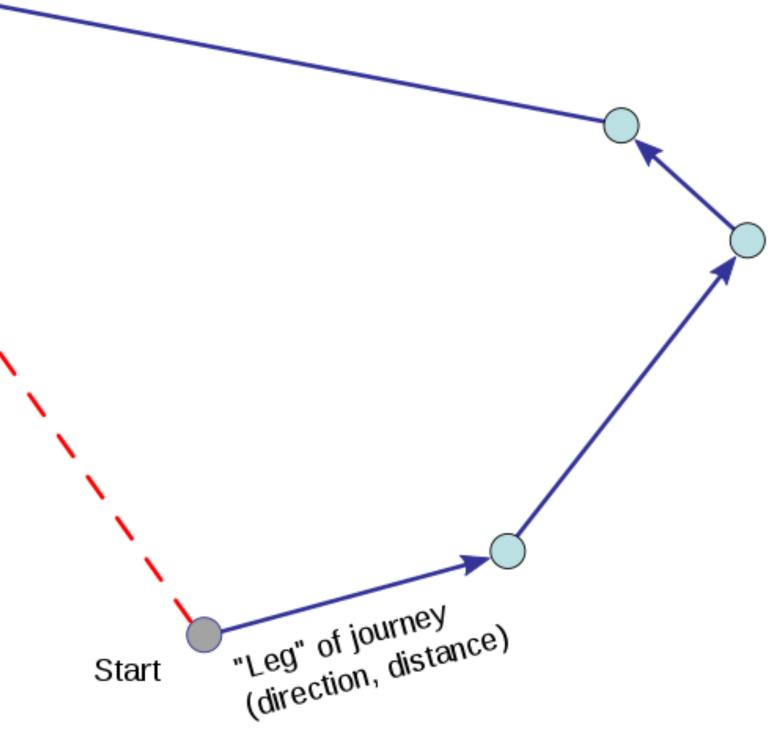
- Basis set to represent functions
- Filtering of scales
- Smoothing and interpolation
- Spectral coordinates
- •

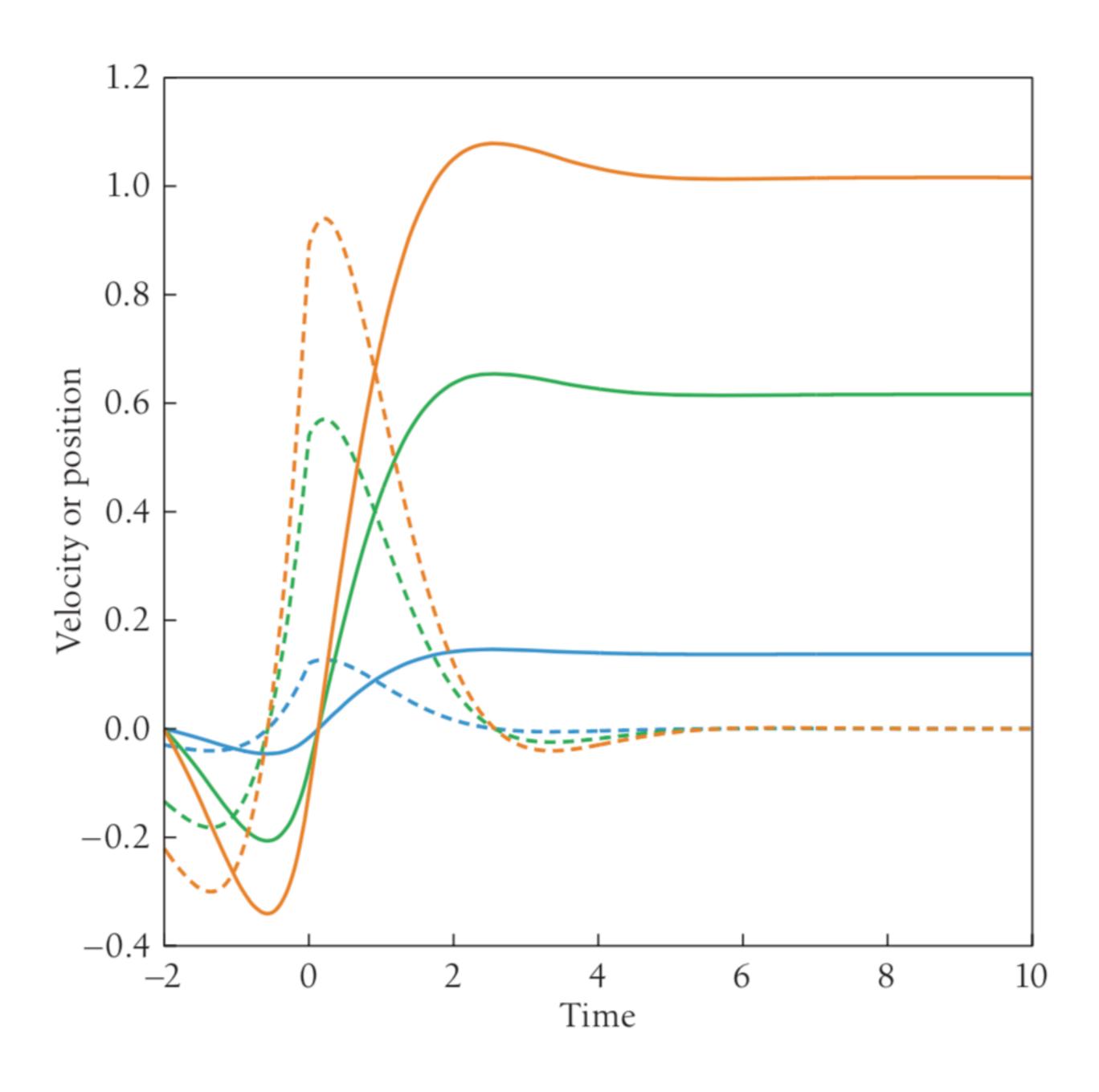
Path integration

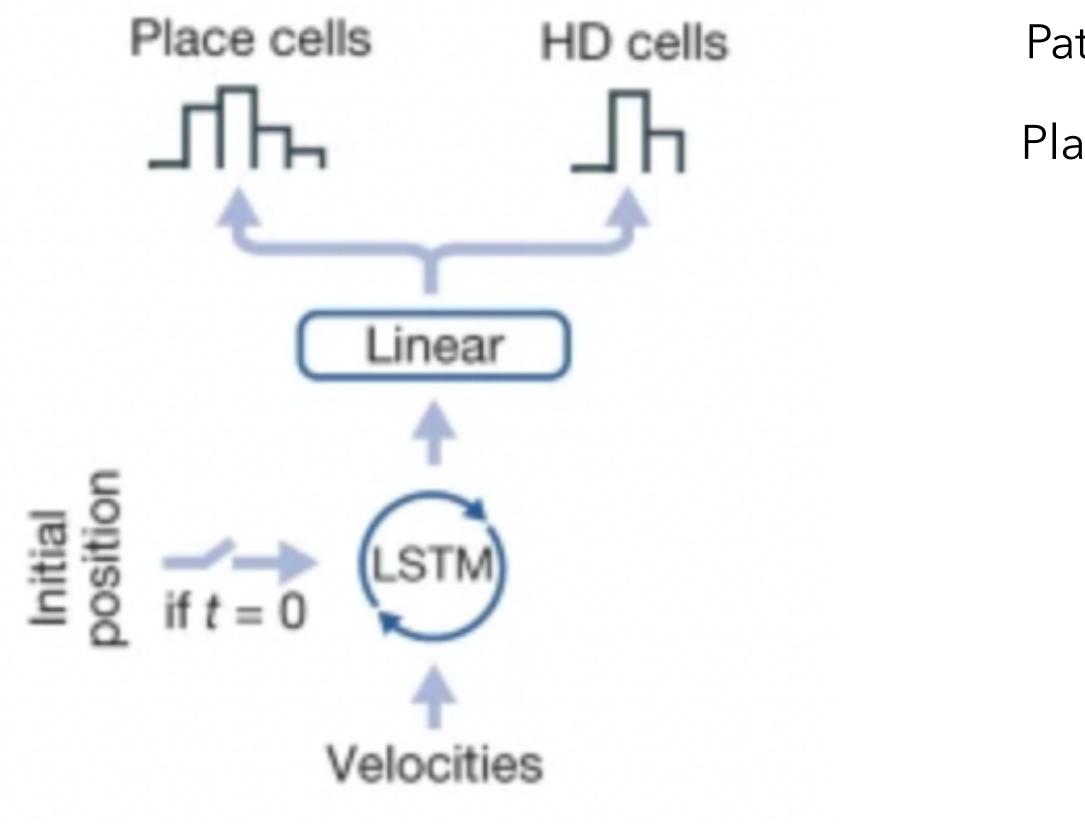


Current Position

Integrated Path estimates Current Position, and gives direction, distance for return journey







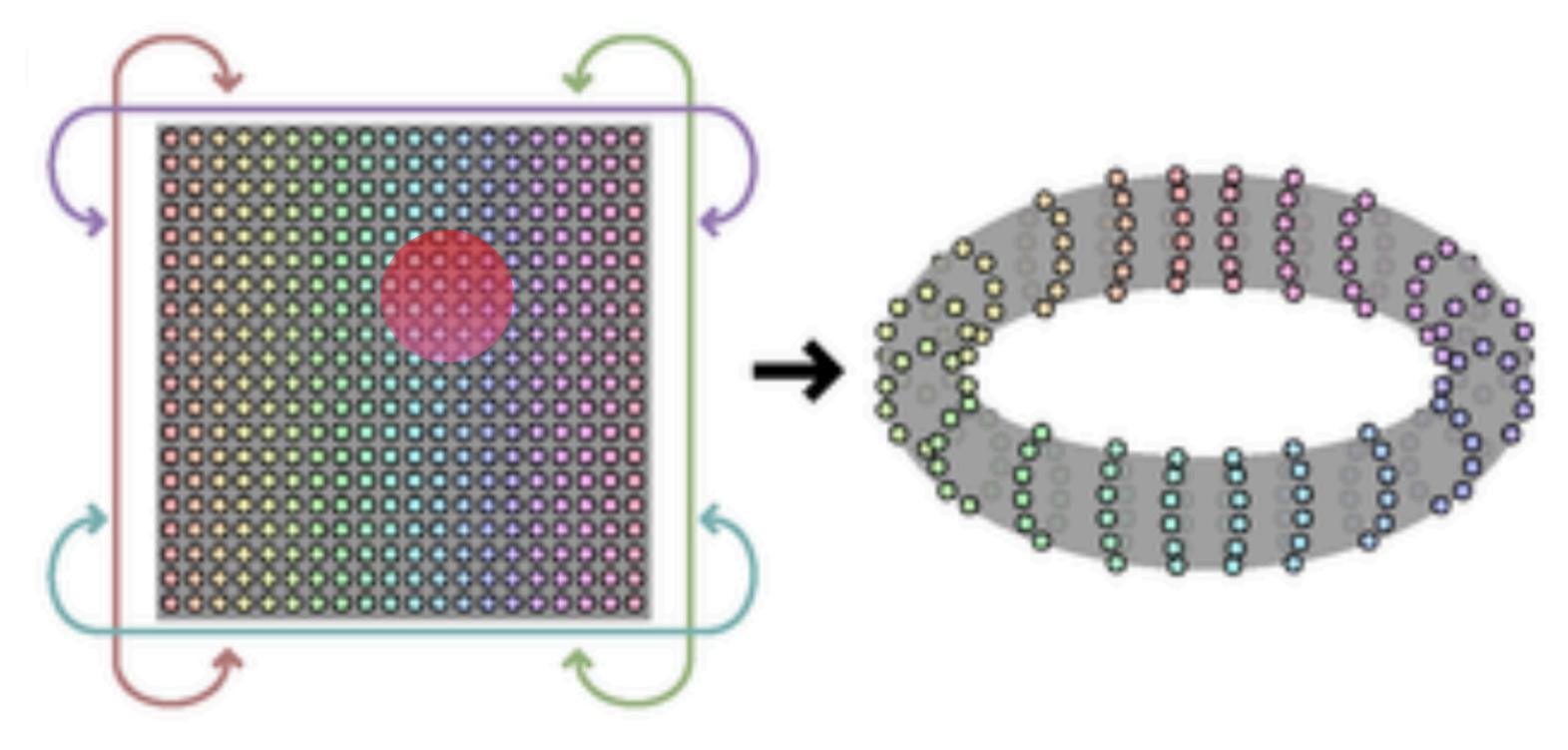
Path integration amounts to compute the integral of **velocity**

Place cells as one-hot encoding of position

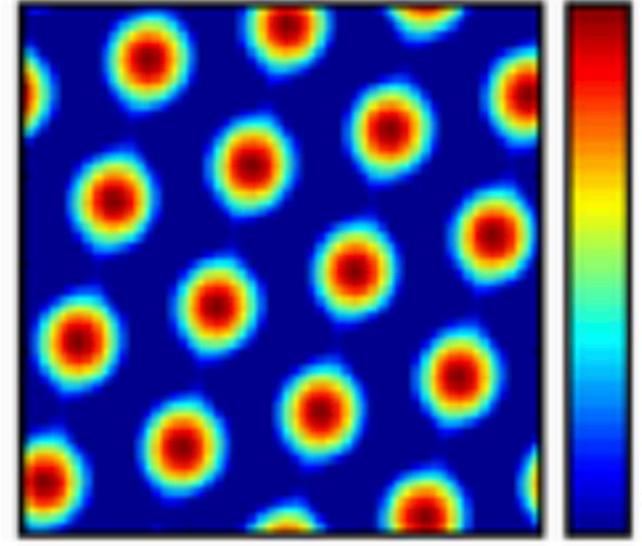


Classical theory of grid cells

- Toroidal connectivity of neurons
- Bump of activity updated by velocity cues (to solve path integration)



Neuronal space



Physical space



Path integration just a particular case of driven dynamical system?

Path integration just a particular case of driven dynamical system?

ivative Form Integral Form r(t) $r(t) = r_0 + \int_0^t v dt'$ Derivative Form

Position

Position and Velocity

$$v(t) = \int \alpha(t) dt + C$$

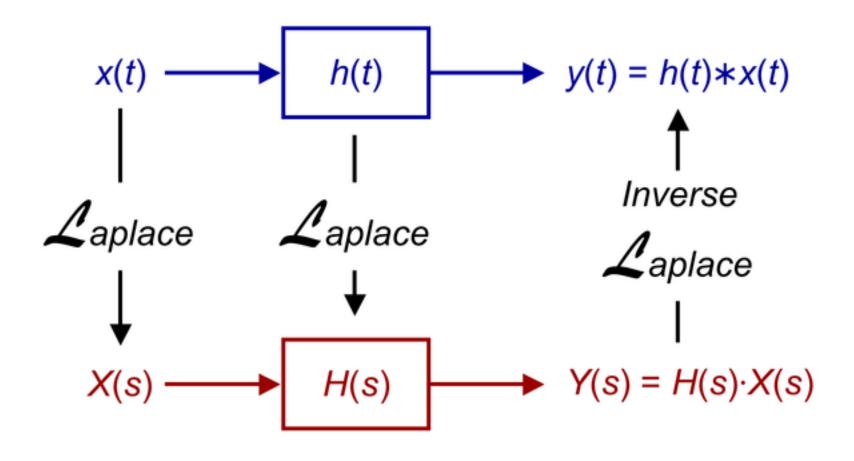
 $\chi(t) = \int v(t) dt + C$

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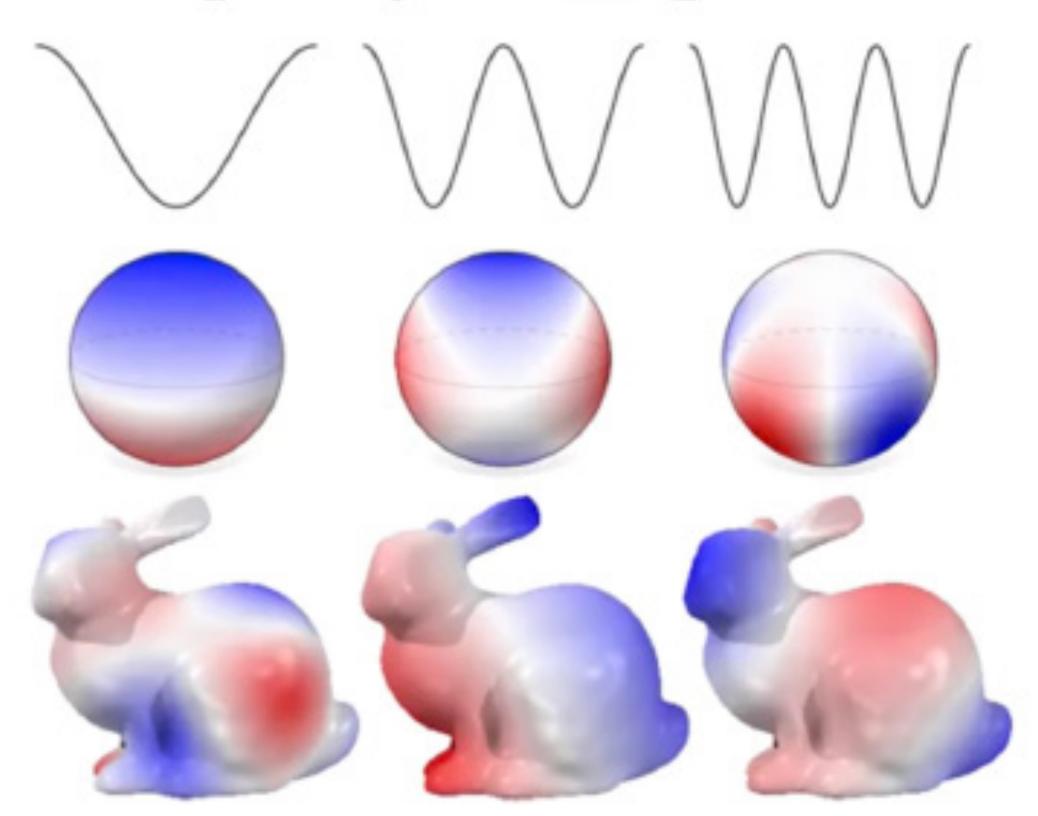
Time domain



Frequency domain

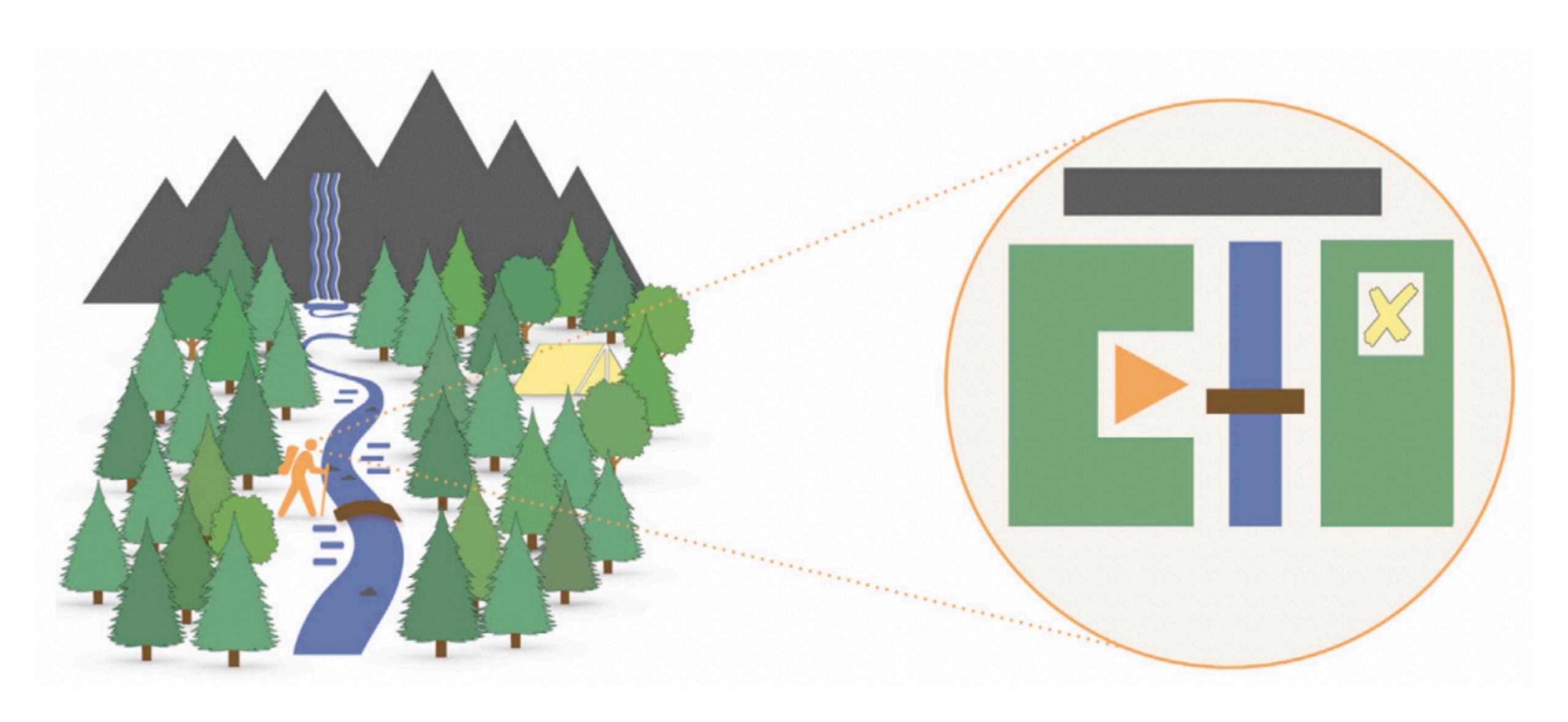
Generalised Fourier

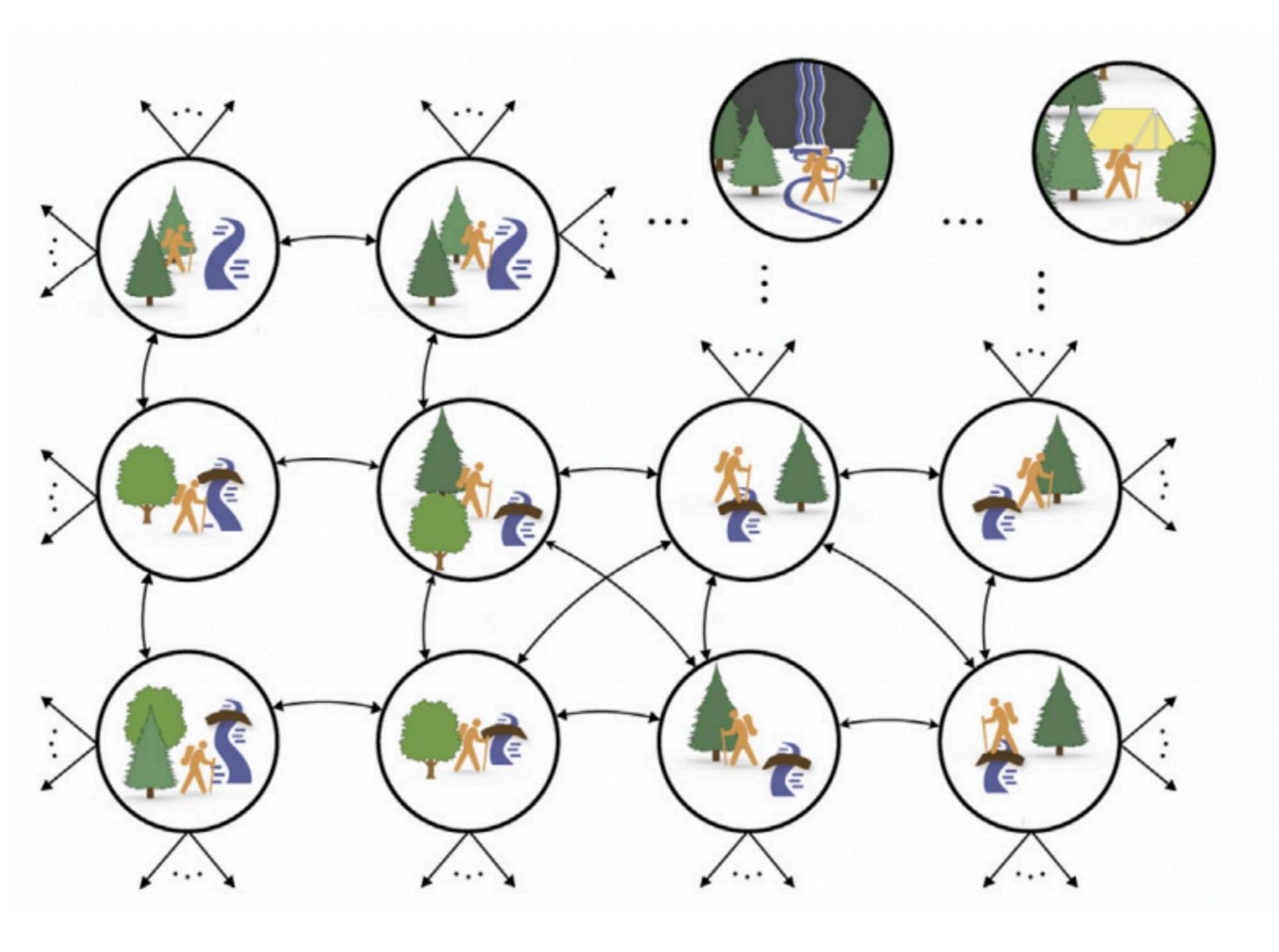
frequency decomposition

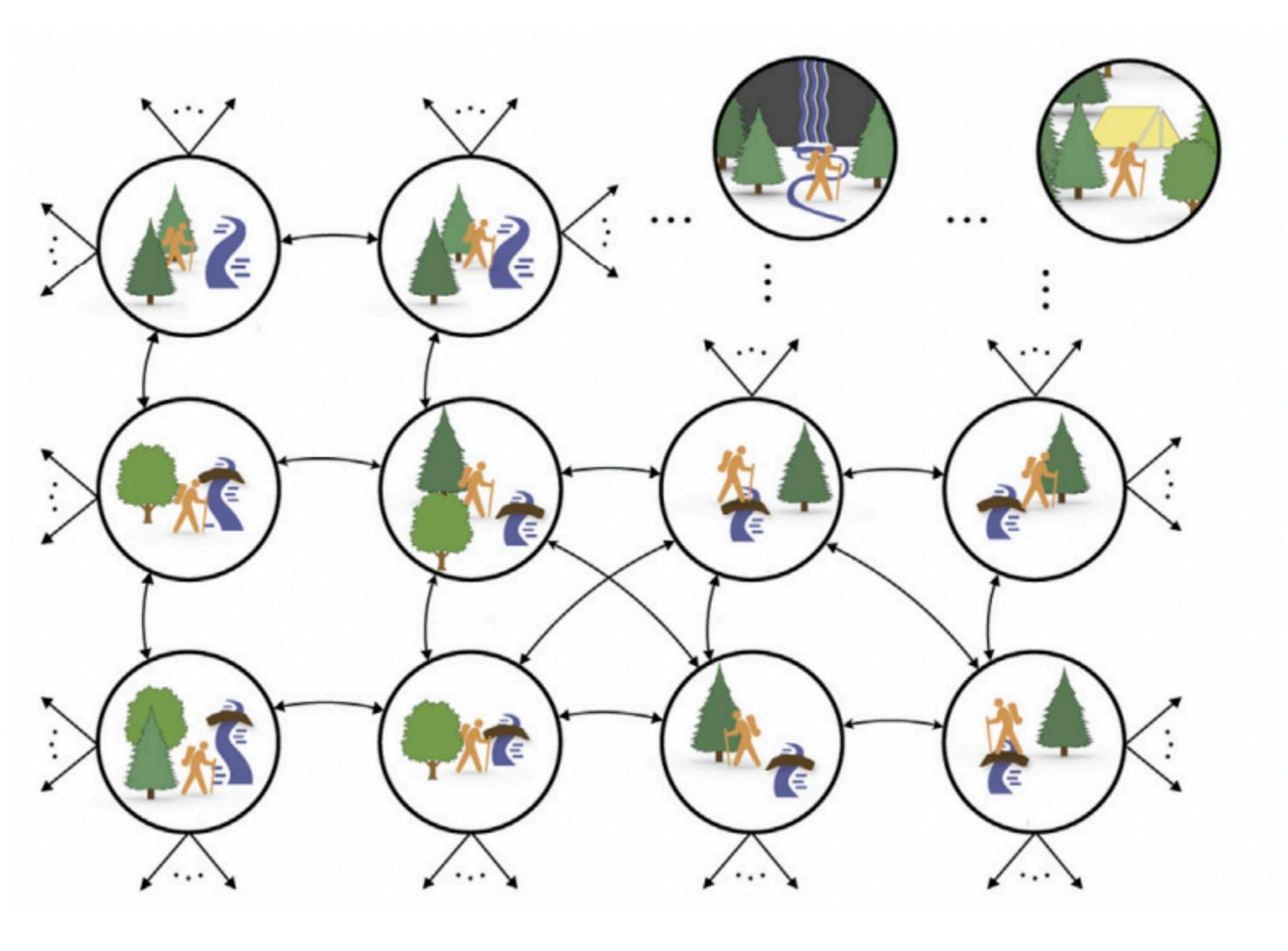


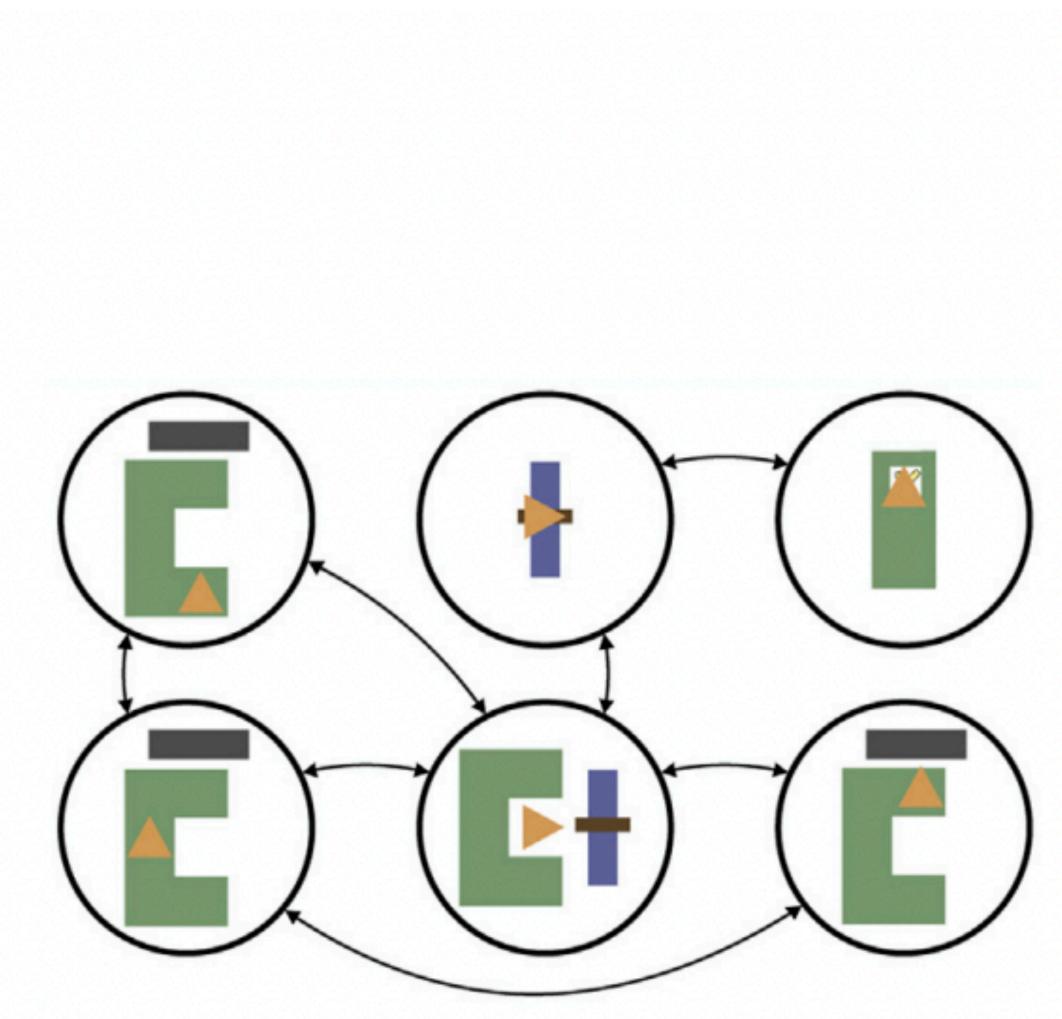
 $\Delta \phi = \lambda \phi$

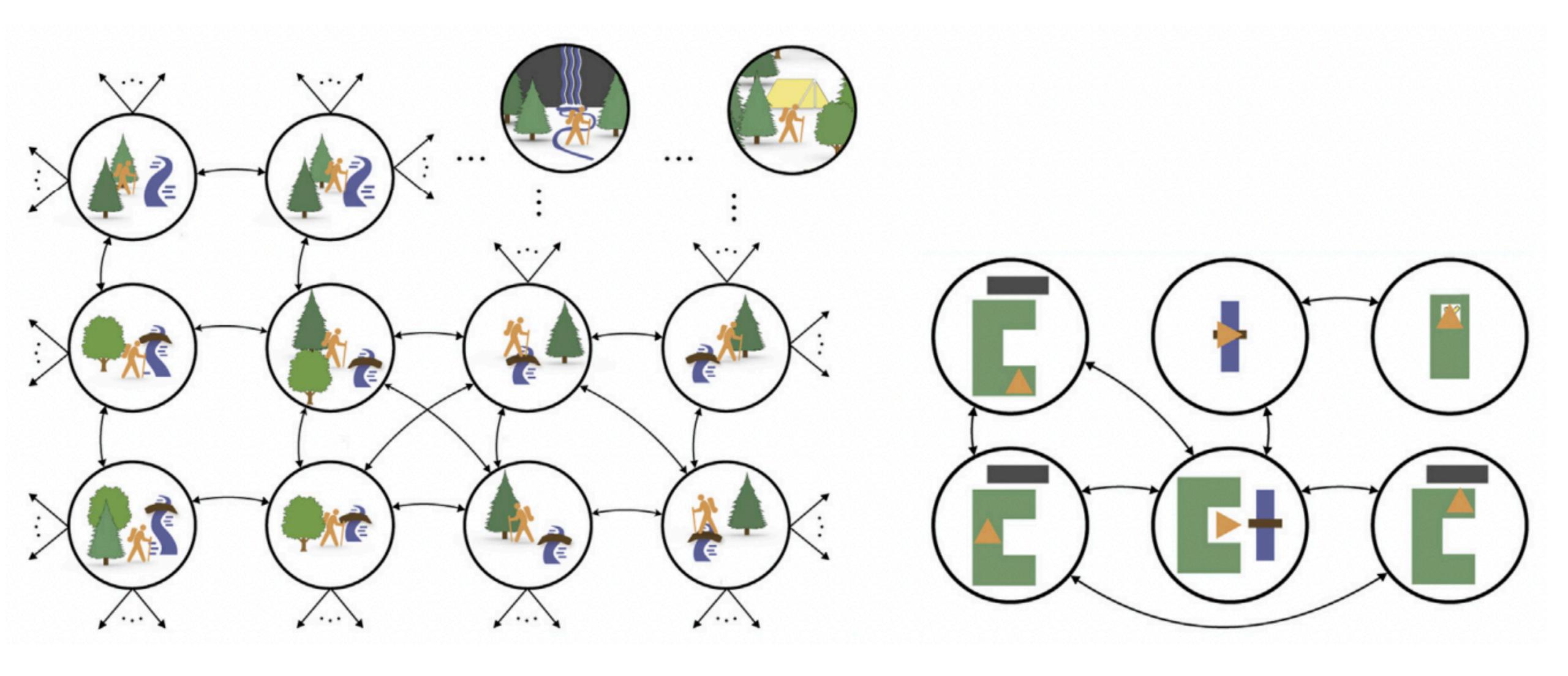
- Basis set to represent functions
- Filtering of scales
- Smoothing and interpolation
- Spectral coordinates
- •



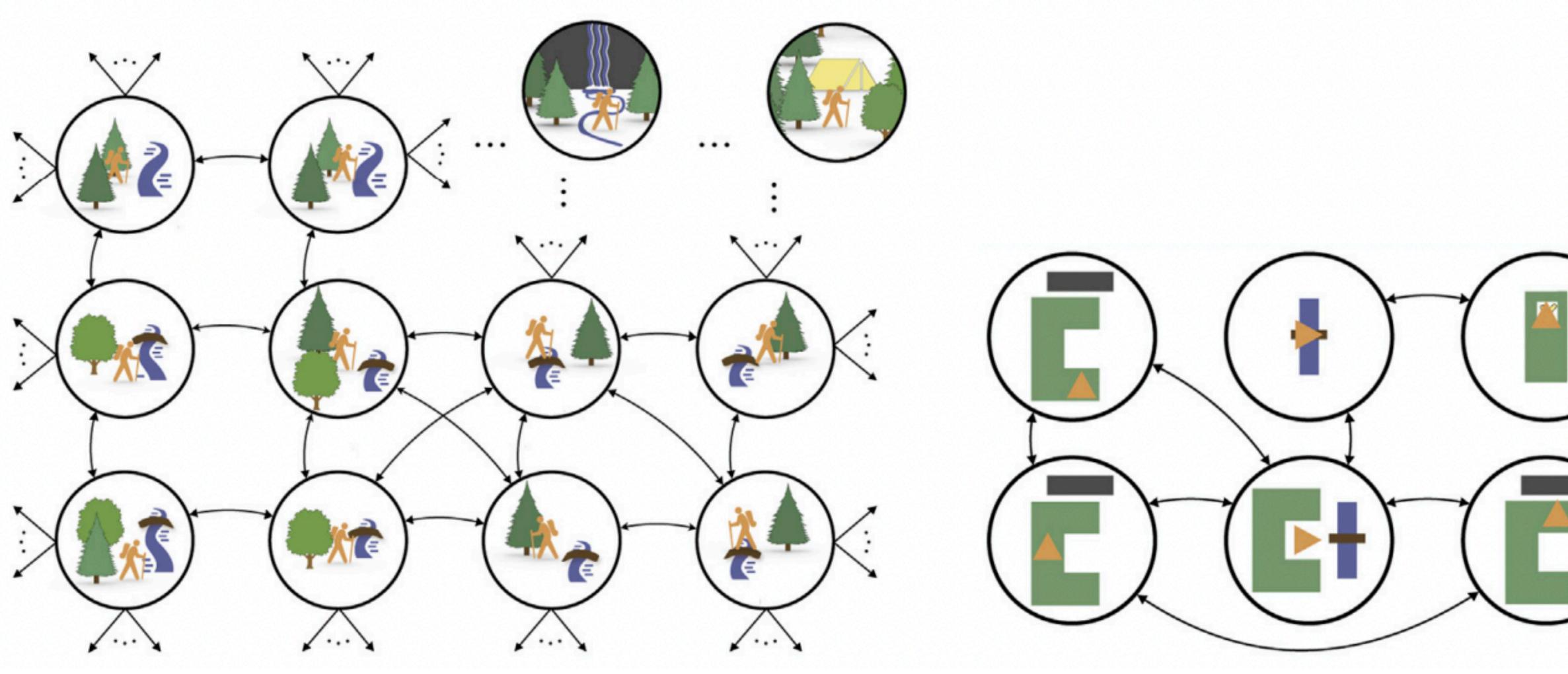








"At the current state, rotate your left leg 25 degrees, place it down beside the rock on the path, then swing your arm forward..."

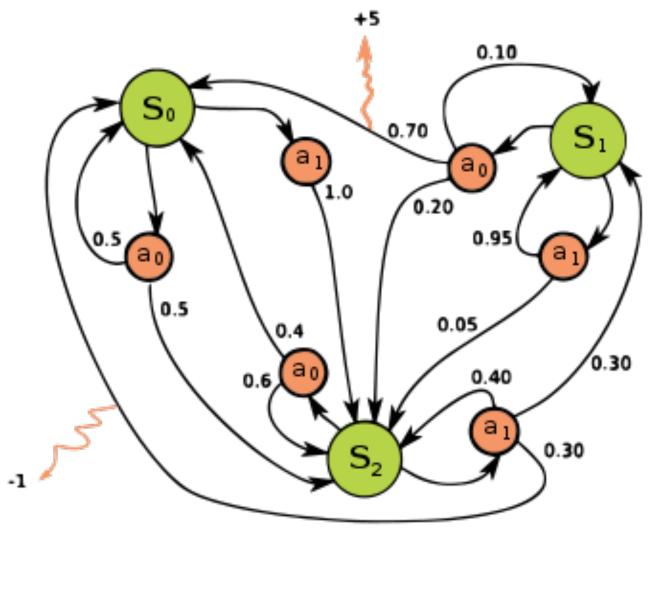


"At the current state, rotate your left leg 25 degrees, place it down beside the rock on the path, then swing your arm forward..." "Since I am on the west side of the river, I probably want to cross the river to get to the east side of the river and then walk towards my campsite."



Reinforcement Learning

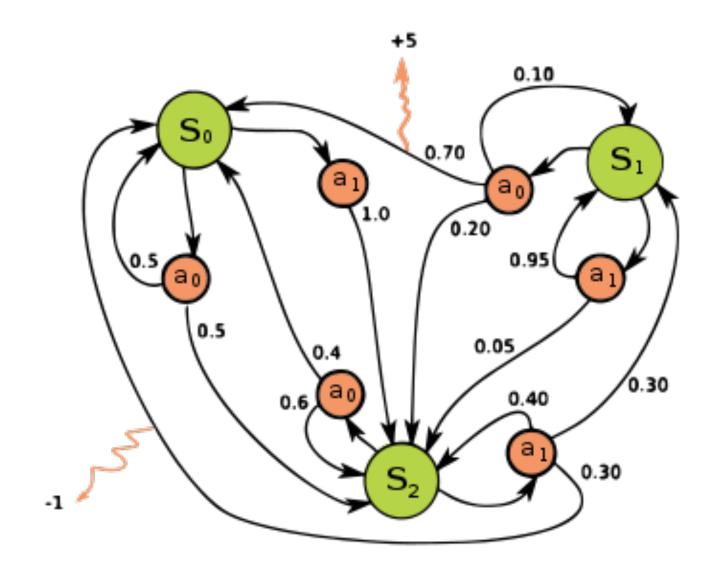
- RL: process to learn long-term future rewards
- Core problem: interdependency of optimal actions
- Tractable when representations encode structure of the environment/task that allows efficient discovery of optimal policy, search of short-paths, replanning, etc...

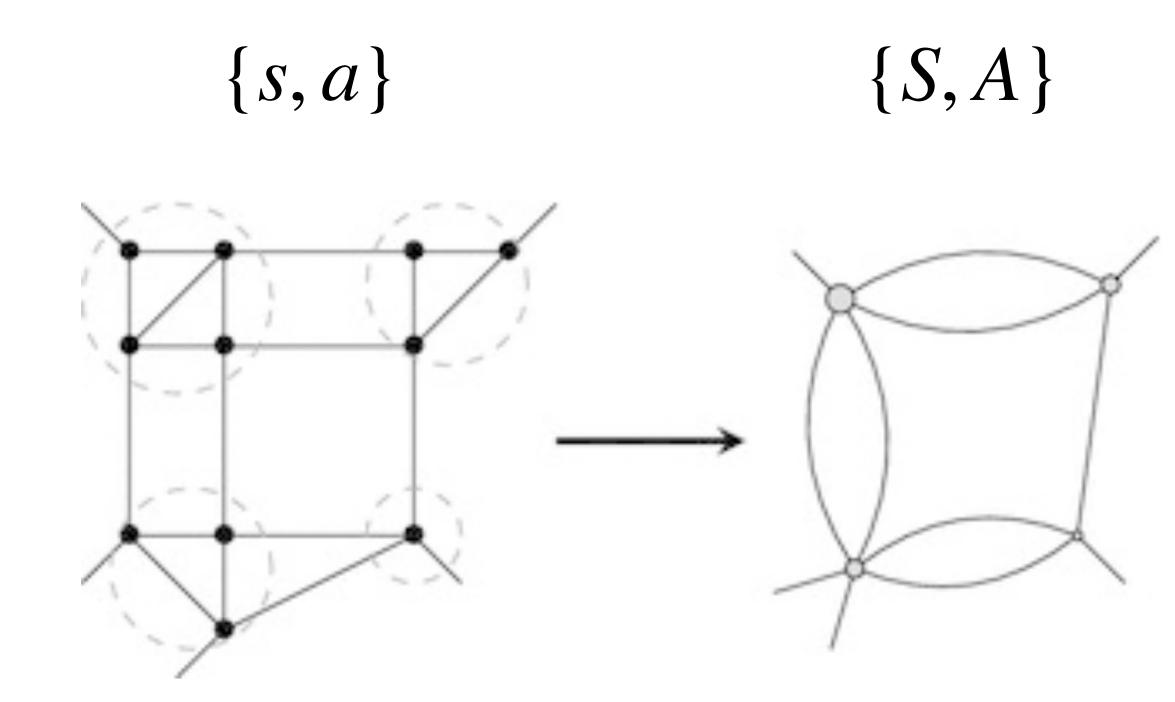


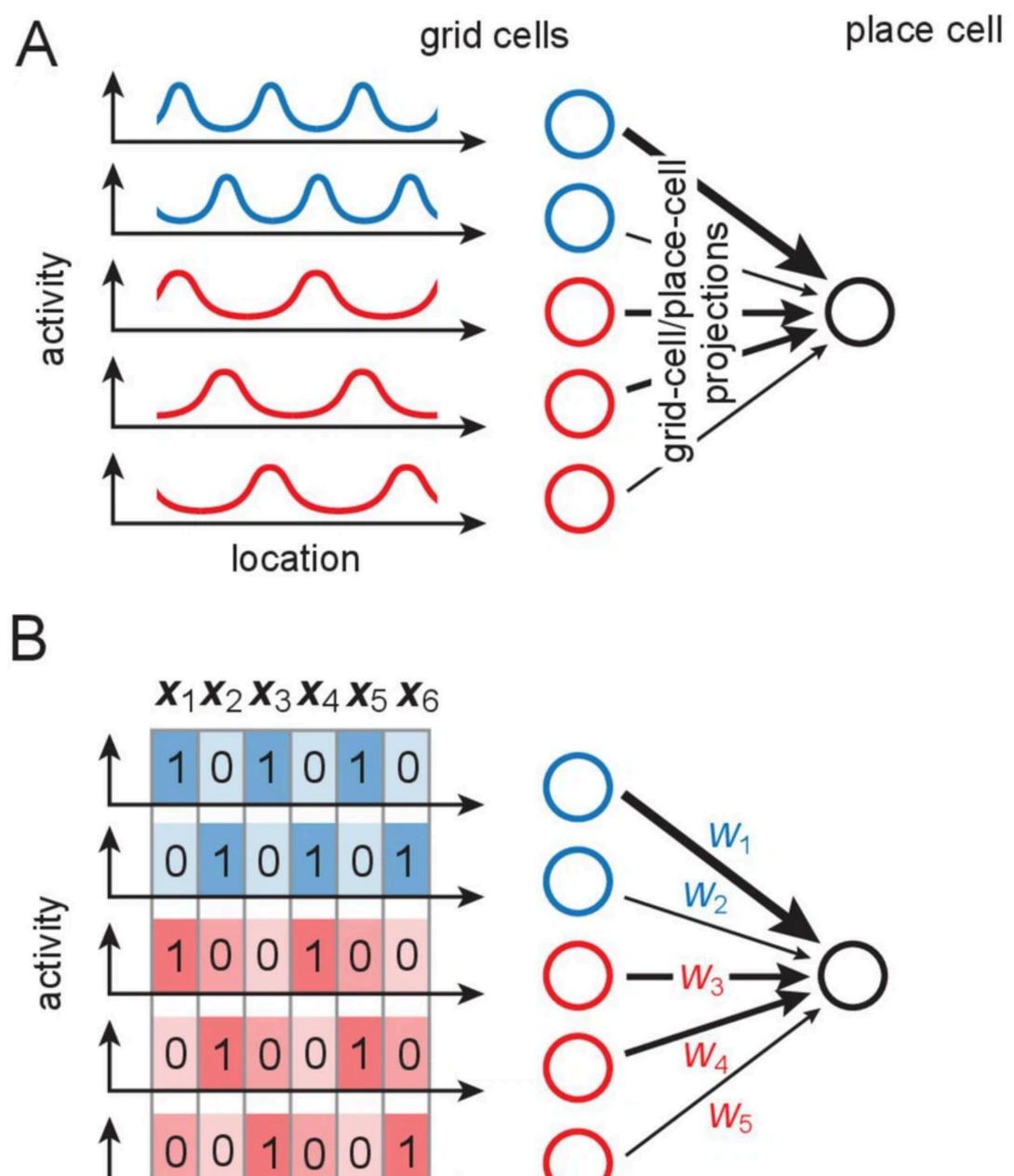
Markov Decision Process

State abstractions: treats certain configurations of the environment as similar by aggregating them. Ex: "being west of the river".

Temporal abstractions: temporally extended macro-actions that describe a general course of action. Ex: "cross the bridge".







123456 j

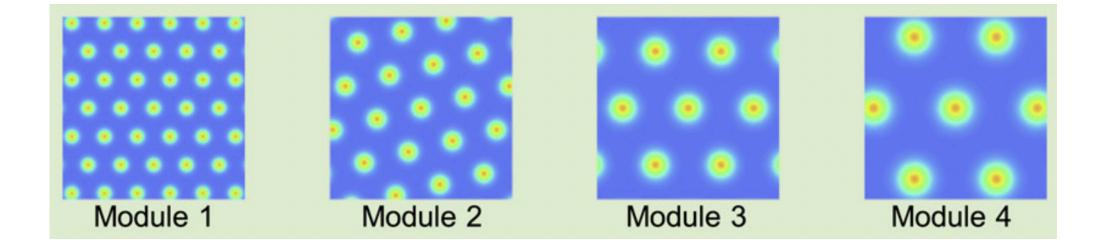
Place cells -----One-hot encoding →

Grid cells – Spectral encoding

 $s(x) = (arphi_1(x), arphi_2(x), \dots, arphi_N(x)) ext{ for vertex } x$

Grid cells

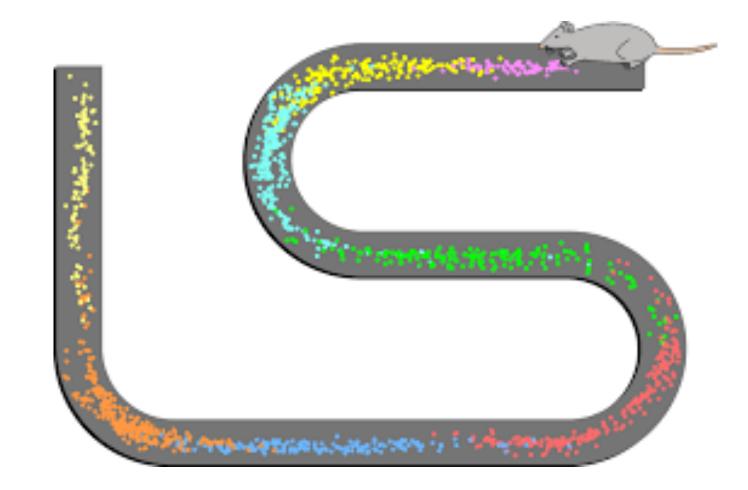
Distributed encoding

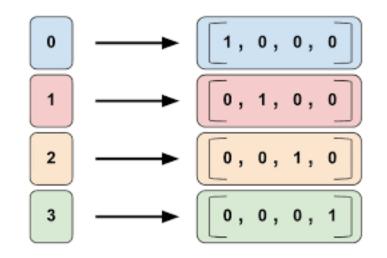


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Place cells

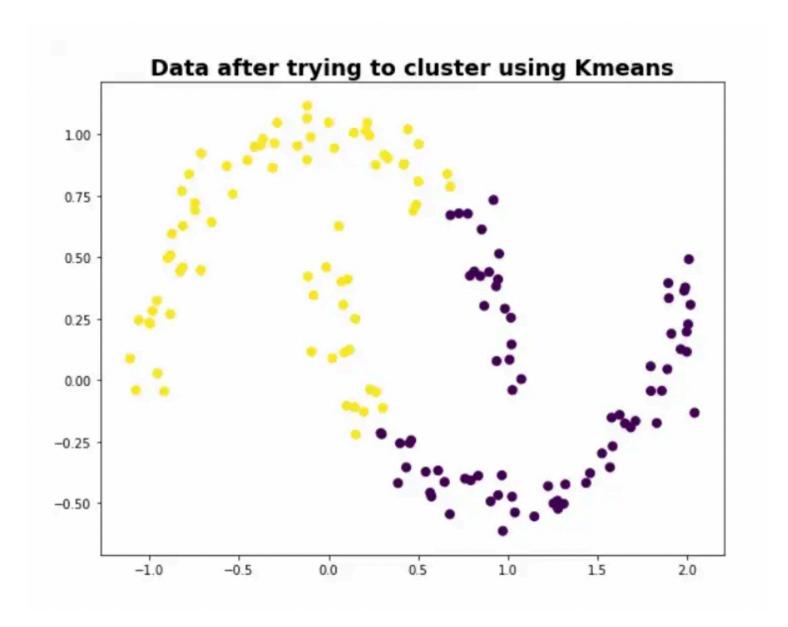
One-hot encoding

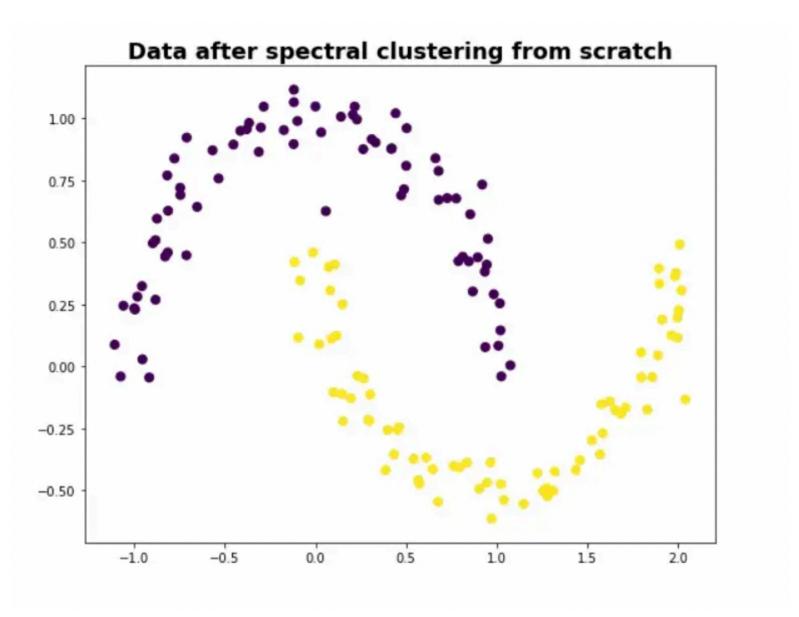




Spectral clustering

- 1. Calculate the Laplacian L (or the normalized Laplacian)
- 2. Calculate the first k eigenvectors (the eigenvectors corresponding to the k smallest eigenvalues of L)
- 3. Consider the matrix formed by the first k eigenvectors; the l-th row defines the features of graph node l
- 4. Cluster the graph nodes based on these features (e.g., using k-means clustering)

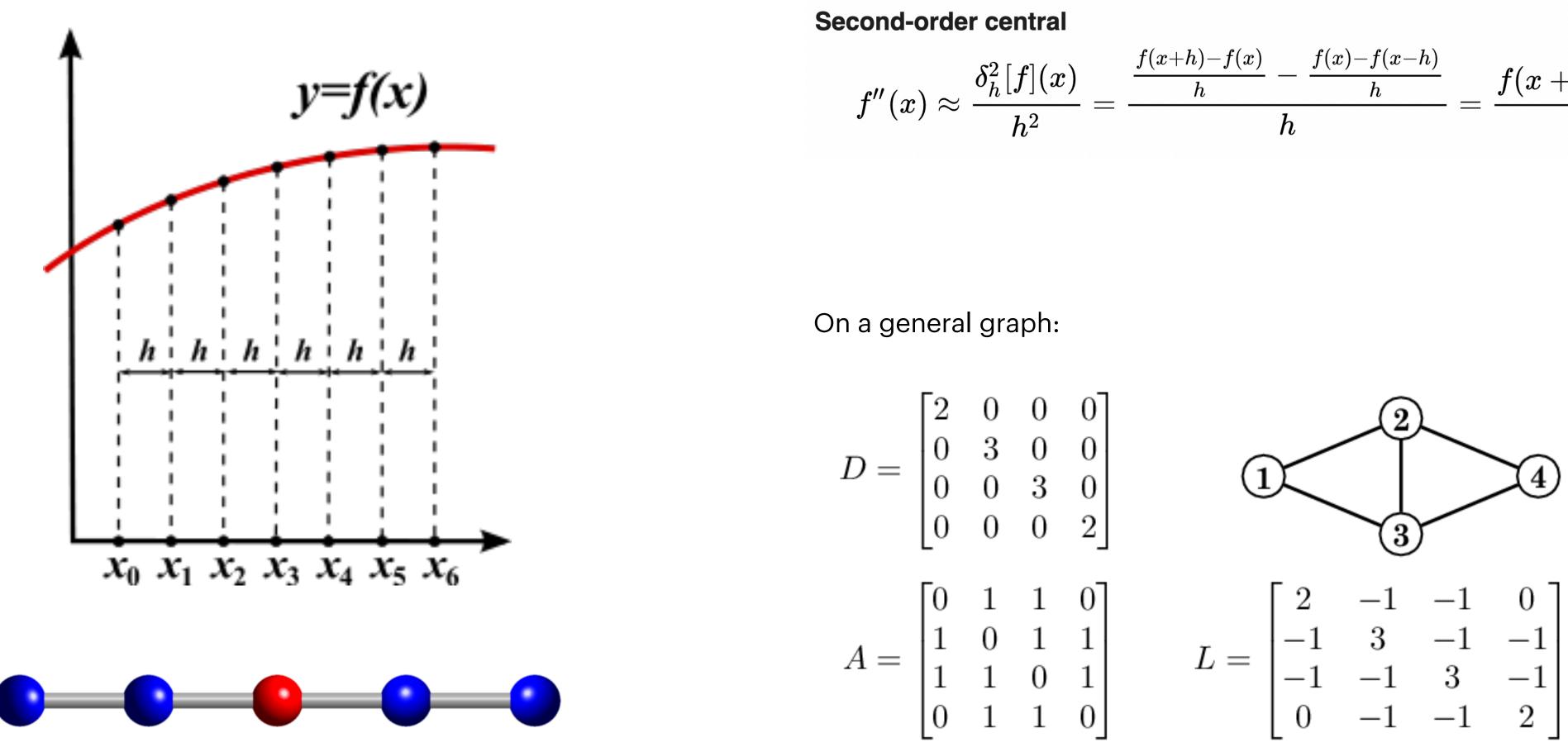




Ng, Andrew Y.; Jordan, Michael I.; Weiss, Yair (2002). "On spectral clustering: analysis and an algorithm" Advances in Neural Information Processing Systems.



Graph Laplacian operator

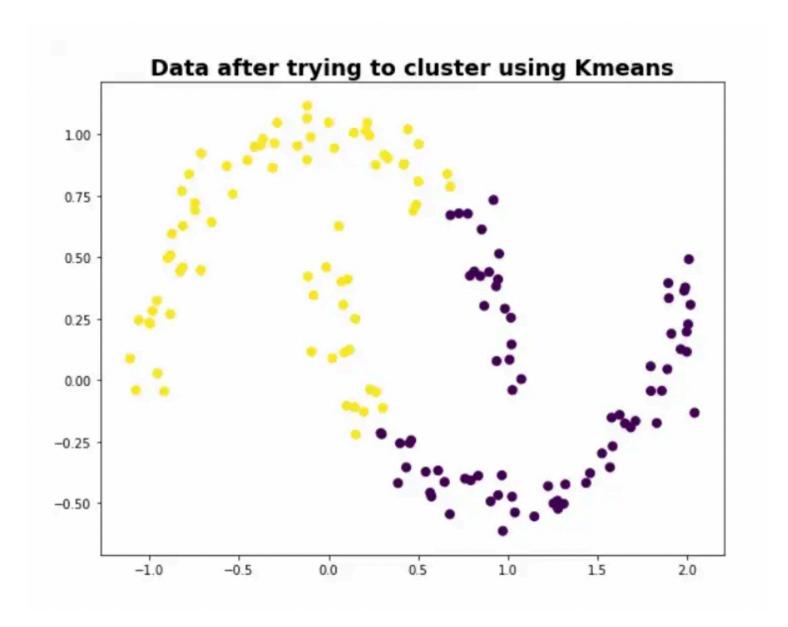


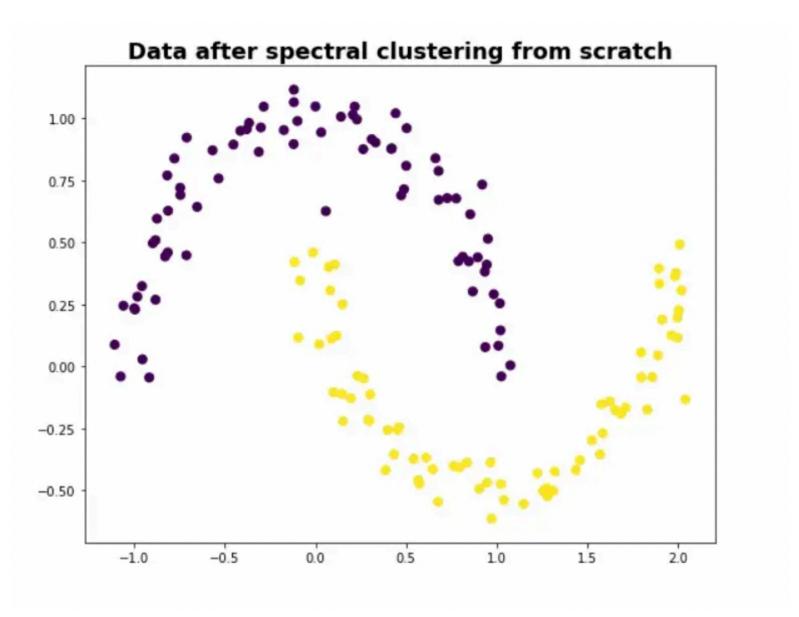
$$\delta_{h}^{2}[f](x) = rac{f(x+h)-f(x)}{h} - rac{f(x)-f(x-h)}{h} = rac{f(x+h)-2f(x)+f(x-h)}{h^{2}}.$$

Kernel (1, -2, 1)

Spectral clustering

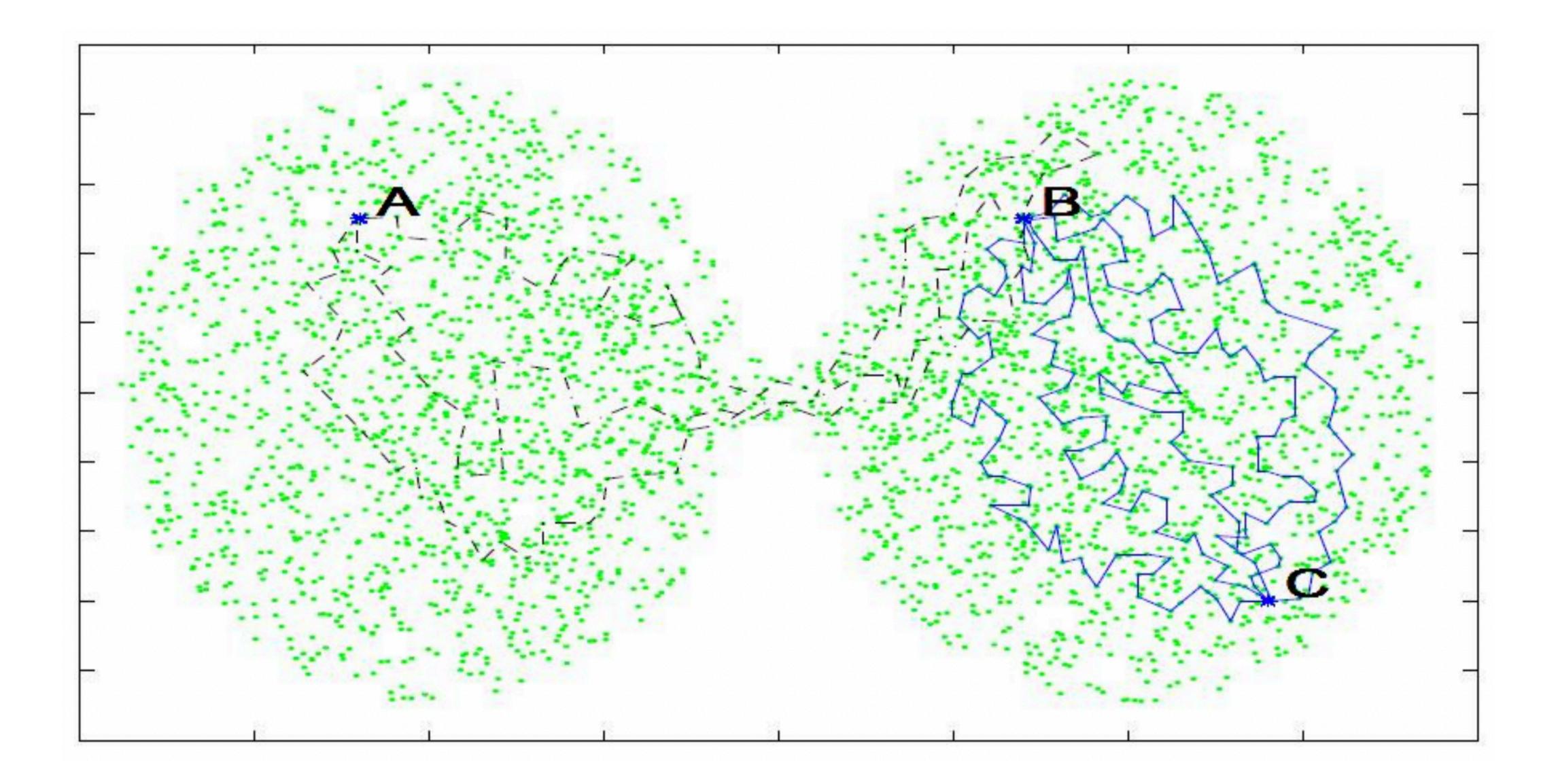
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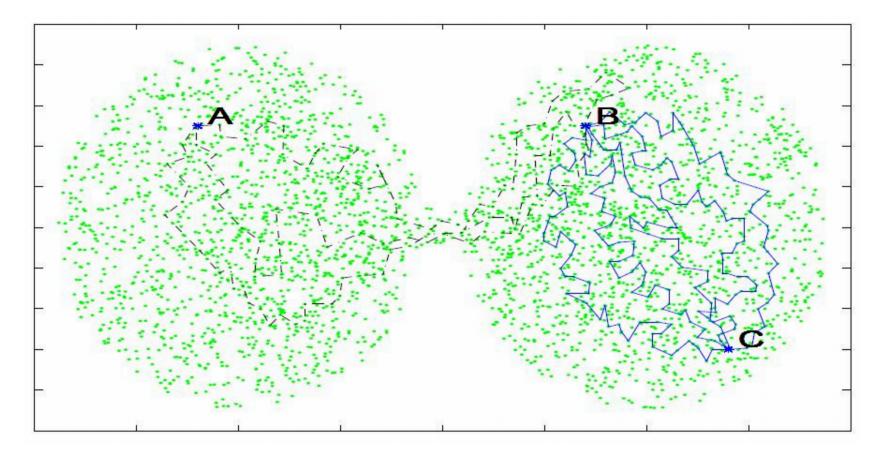


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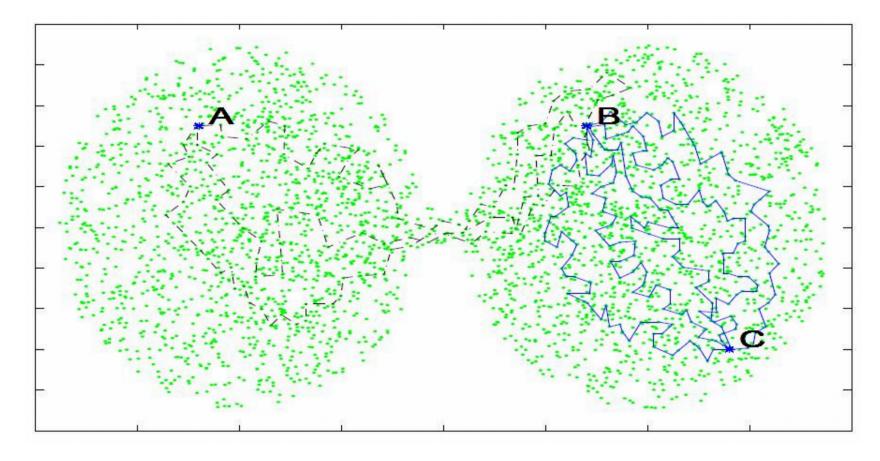


- Dimensionality reduction technique
- Embedding of data set into Euclidean space (low dimension)
- Coordinates computed from eigenvectors and eigenvalues of diffusion operator
- Geometric structure of high-dimensional data by modeling a diffusion process on a graph.



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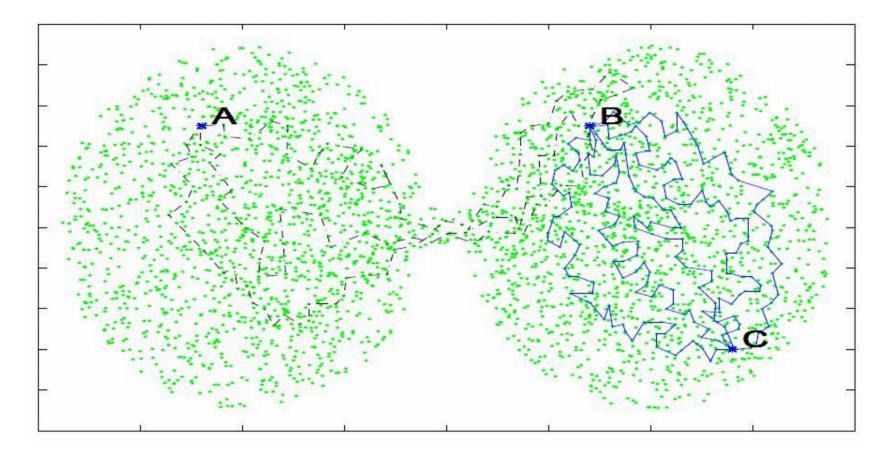
| Matrix | Definition | Purpose in Diffusi |
|---|---|-----------------------------|
| Affinity Matrix W | $W_{ij}=e^{-rac{\ x_i-x_j\ ^2}{\sigma^2}}$ | Defines similarity b |
| Degree Matrix D | $D_{ii} = \sum_j W_{ij}$ | Measures total cor point |
| Laplacian $L=D-W$ | Standard graph Laplacian | Captures local diff |
| Random Walk Laplacian $L_{ m rw}=I-D^{-1}W$ | Scales diffusion process | Used to define Ma |
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| Eigenvectors of P or $L_{ m rw}$ | Basis for diffusion maps | Used for dimension |



| sion Maps |
|---------------------|
| between points |
| onnectivity of each |
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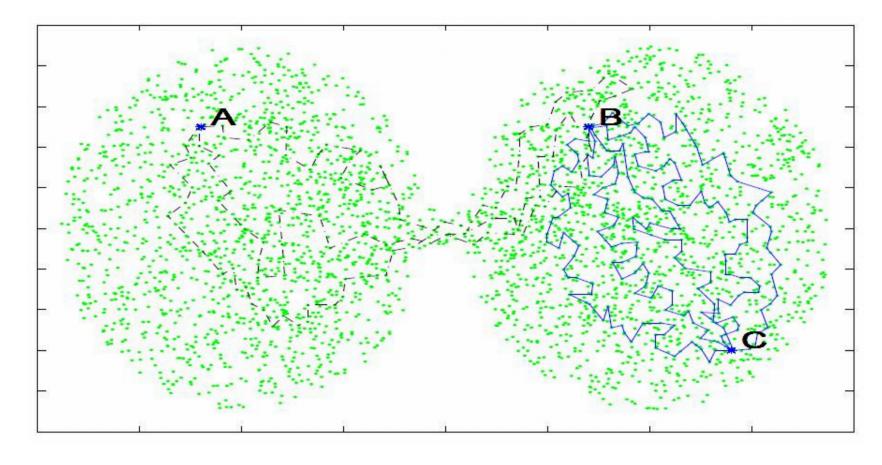
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$(y,t|x_j))^2$

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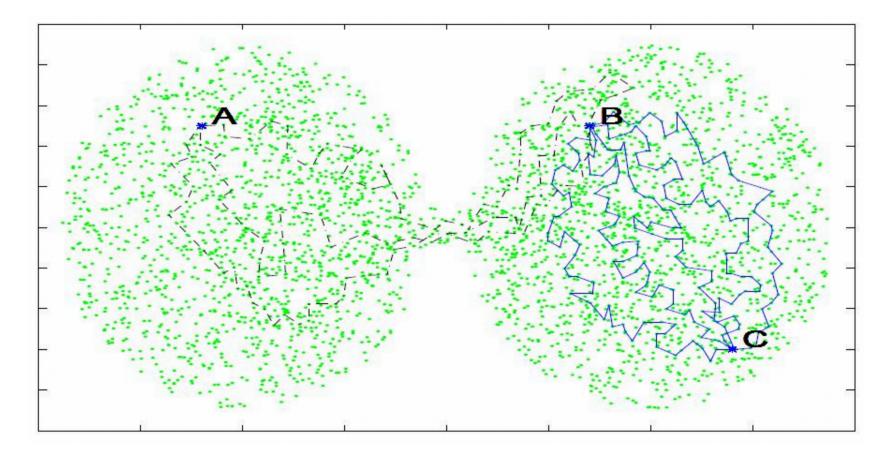
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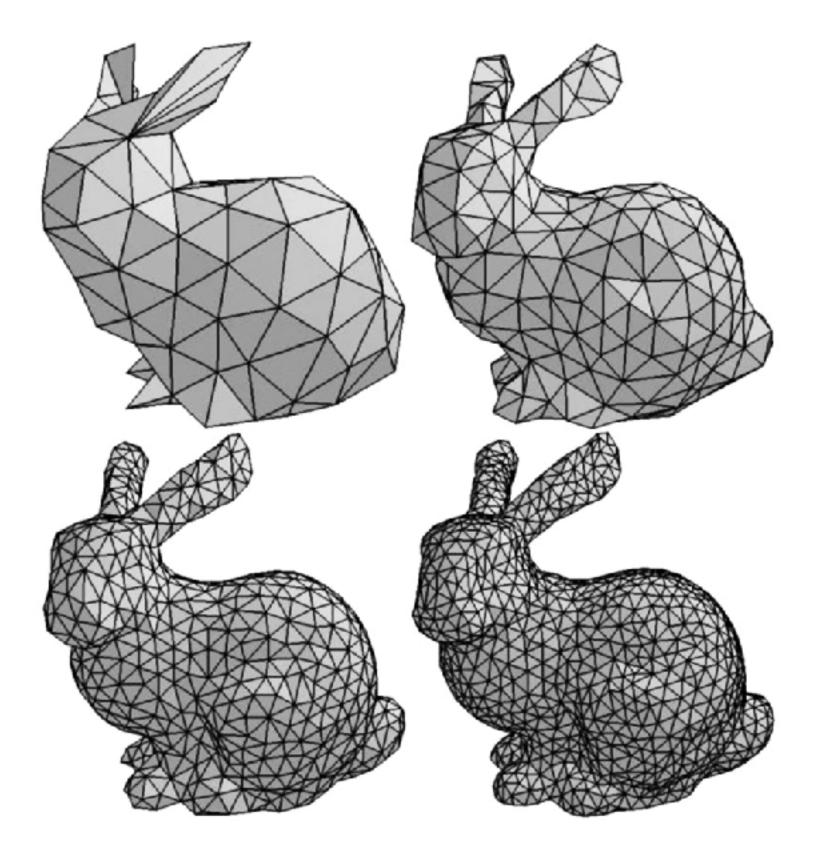
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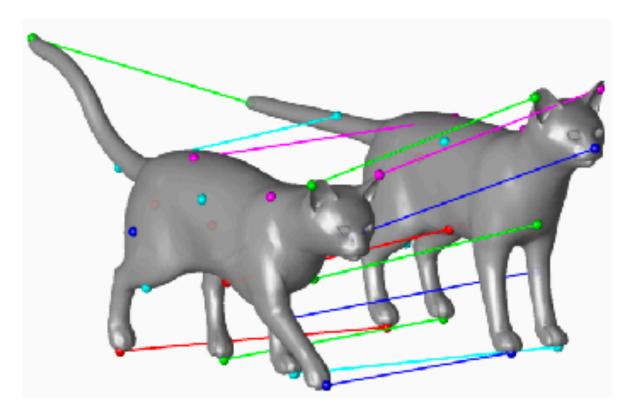
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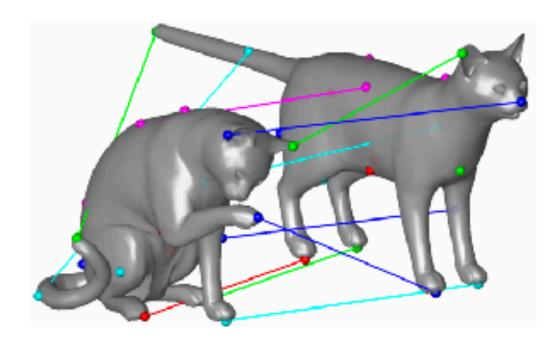


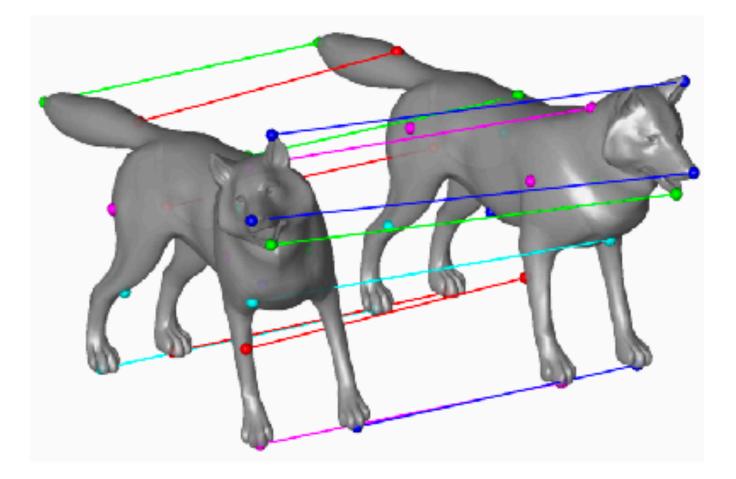


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Spectral shape matching

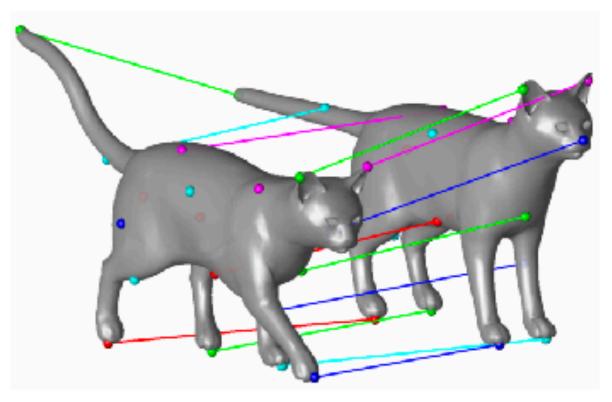


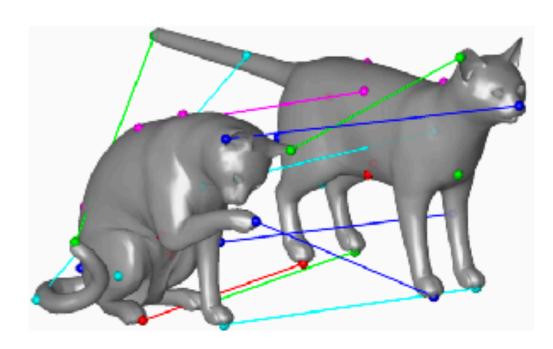


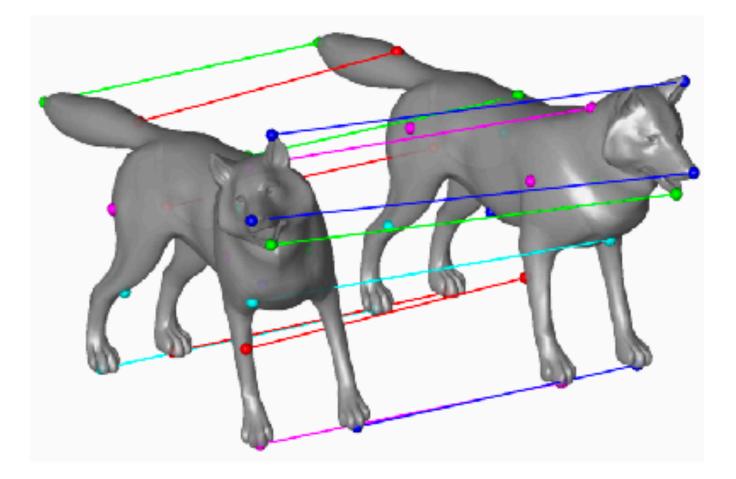


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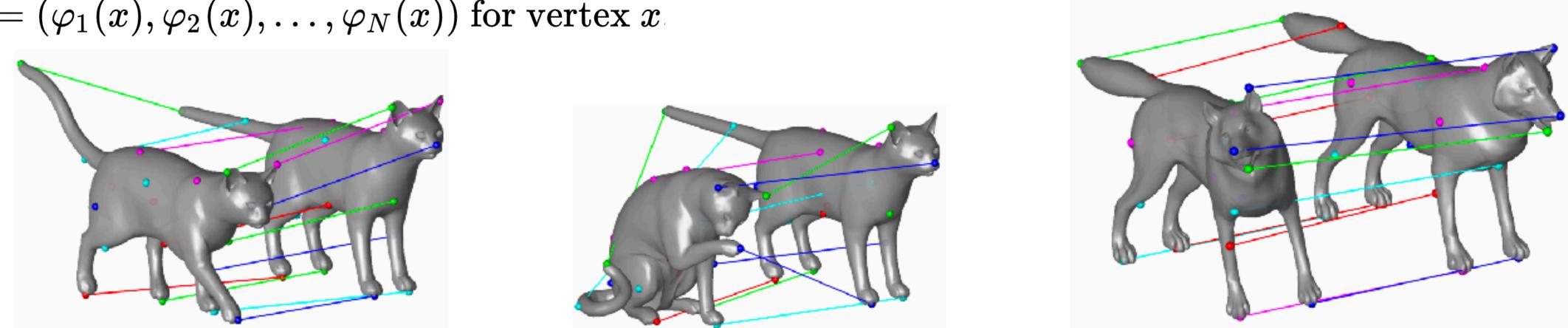




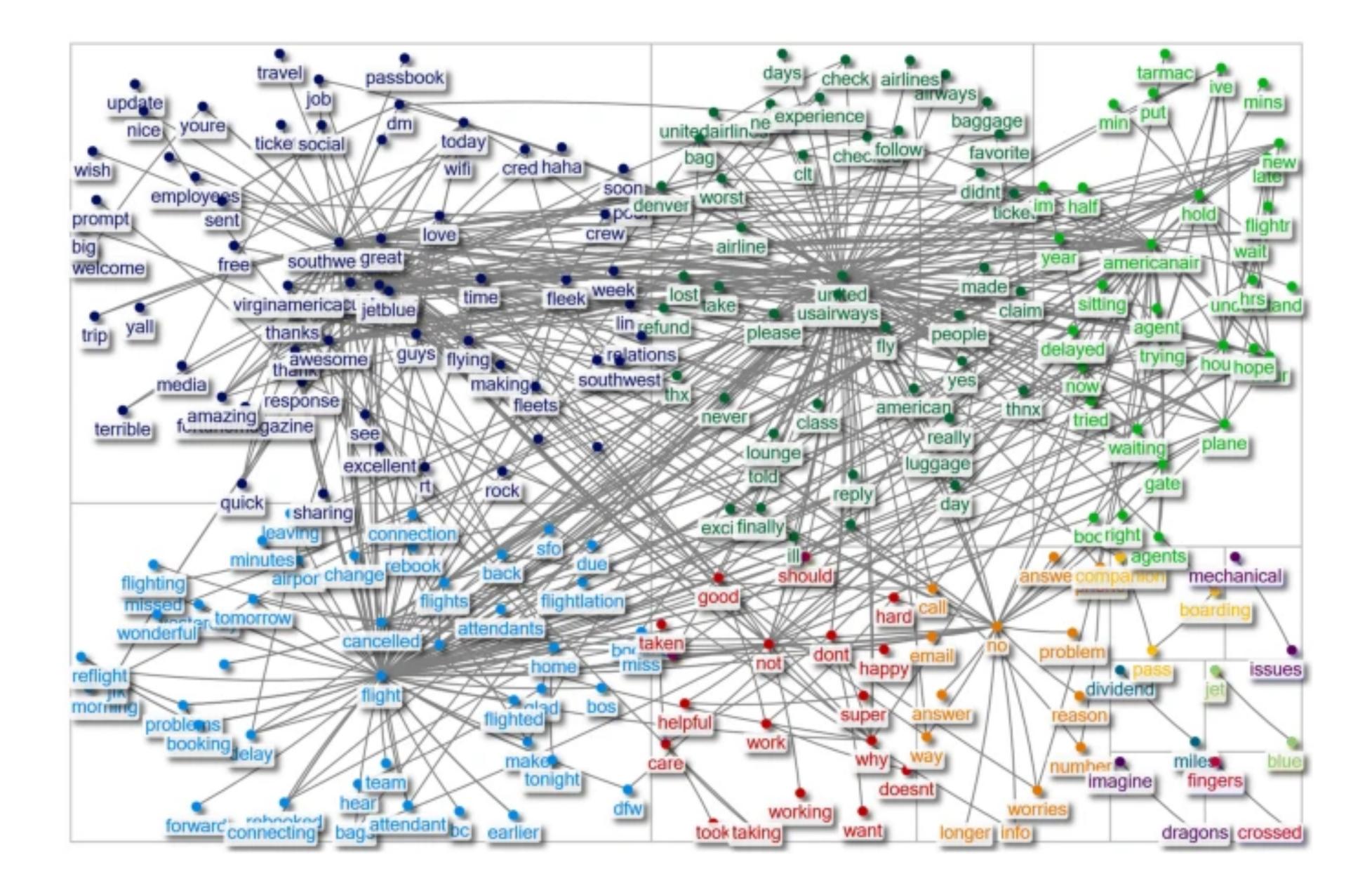


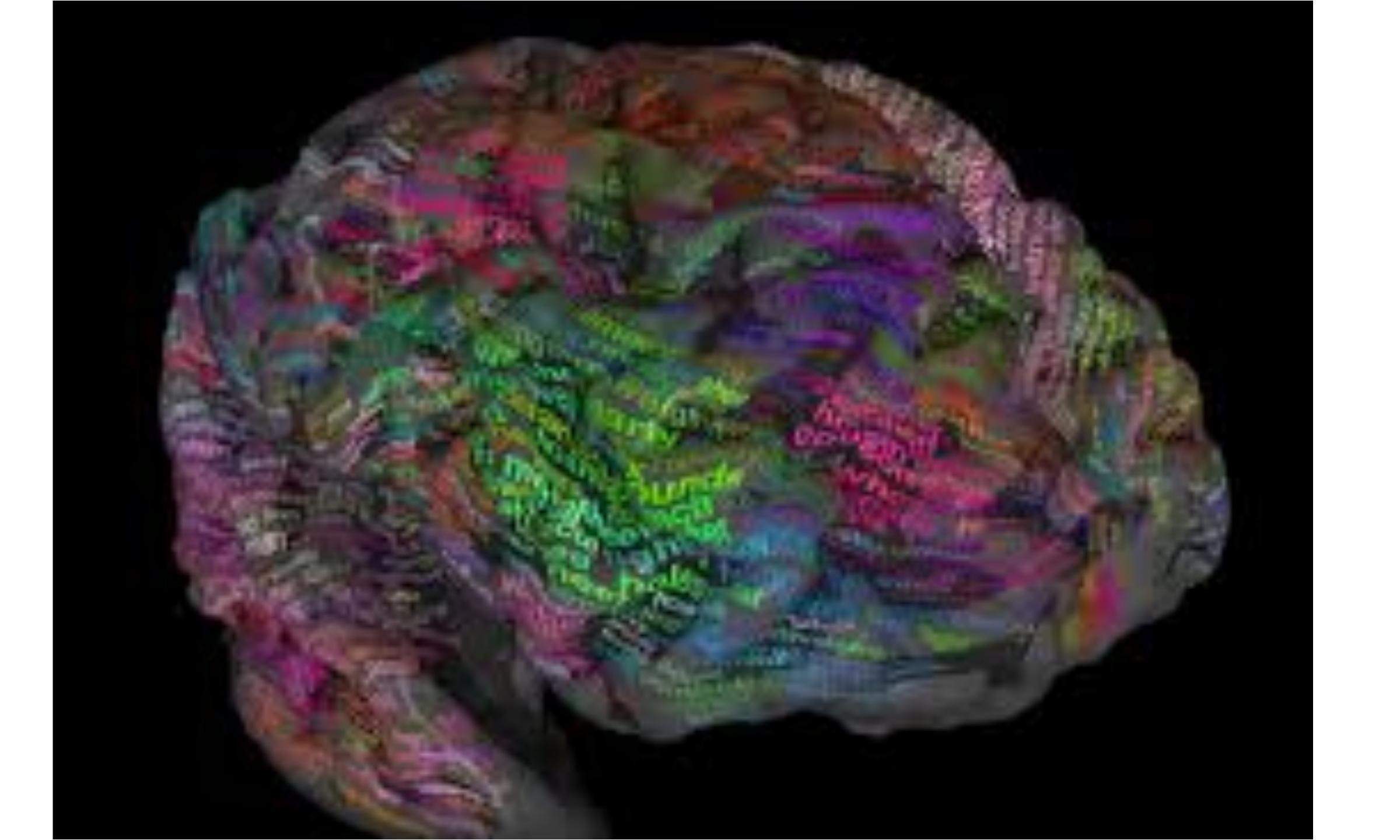
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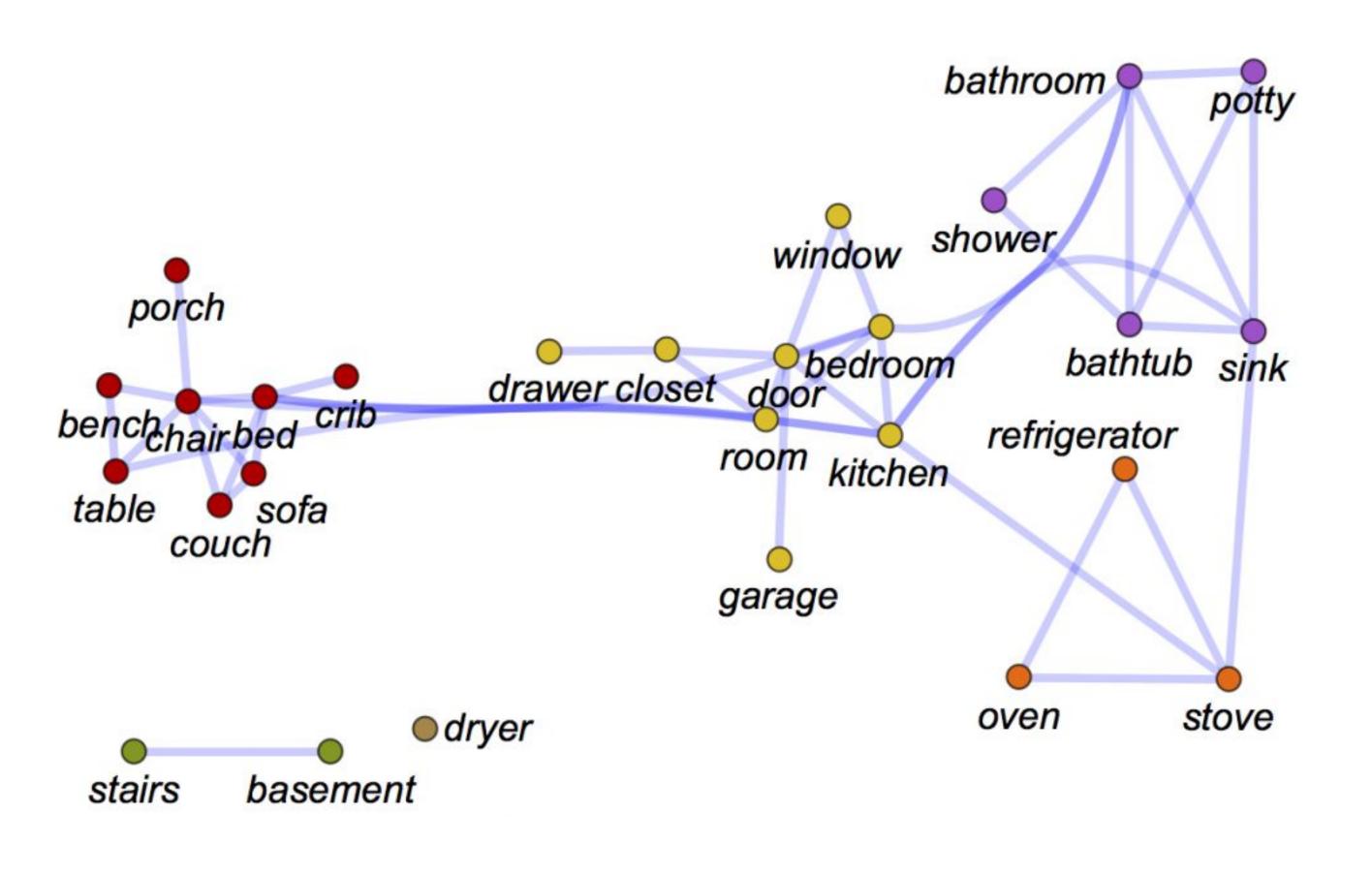


- Low-dimensional representations based on connectivity
- Diffusion distance shows invariance to non-stretching or non-topology changing transformations
- Correspondence problem
- Analogy-based inferences?





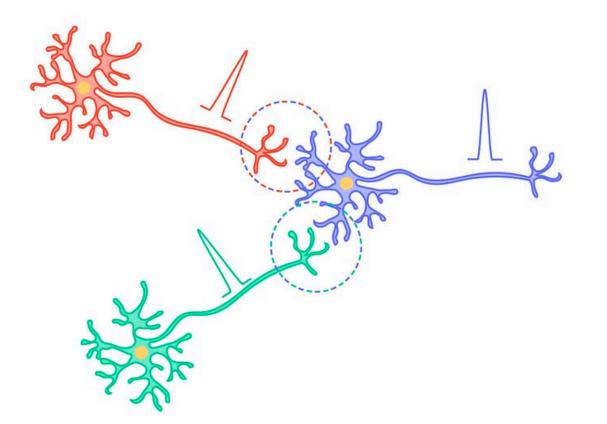
Semantic networks

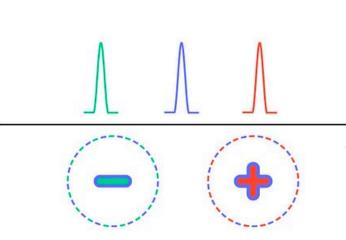


Associative learning

S, S' S, S' S, S'

Temporal sequences s, s', s'', s''', \cdots





 $W_{s,s'}$ \uparrow $W_{s',s}$ \downarrow

Sculpting of neural circuits by weight changes

Hebbian learning (e.g. STDP)

time

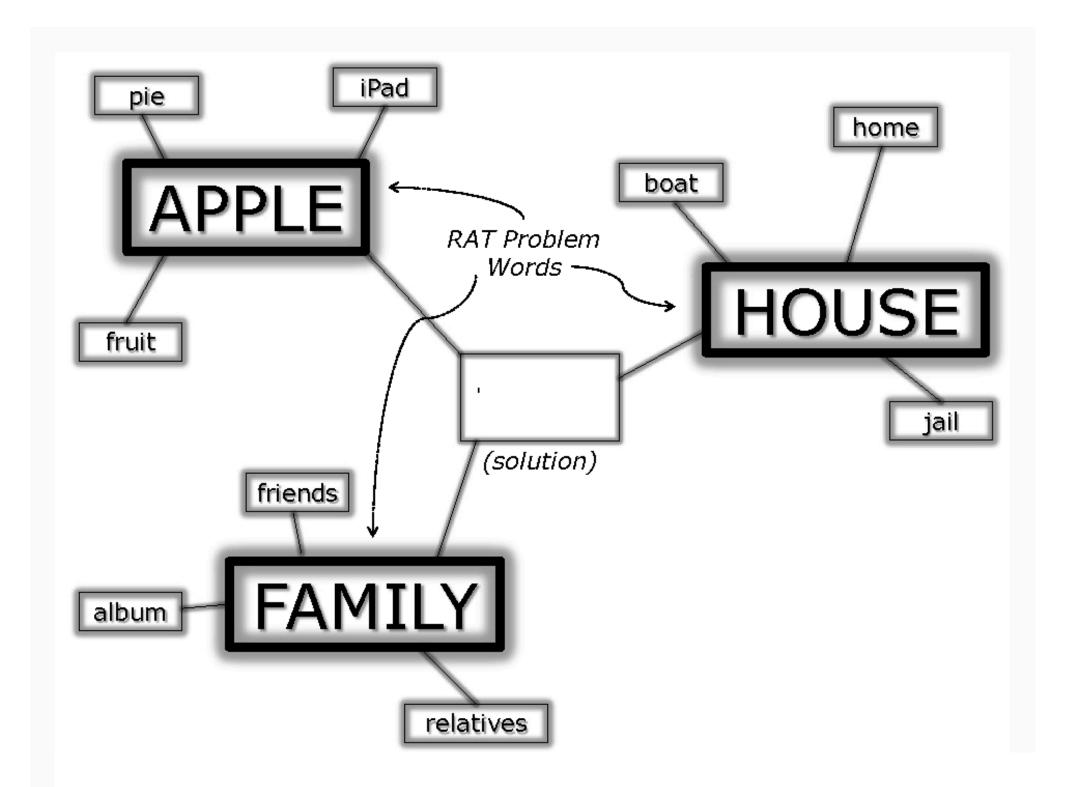
| Word 1 | Word 2 | Word 3 | Answer |
|--------|--------|---------|--------|
| man | glue | star | |
| dew | comb | bee | |
| rain | test | stomach | |

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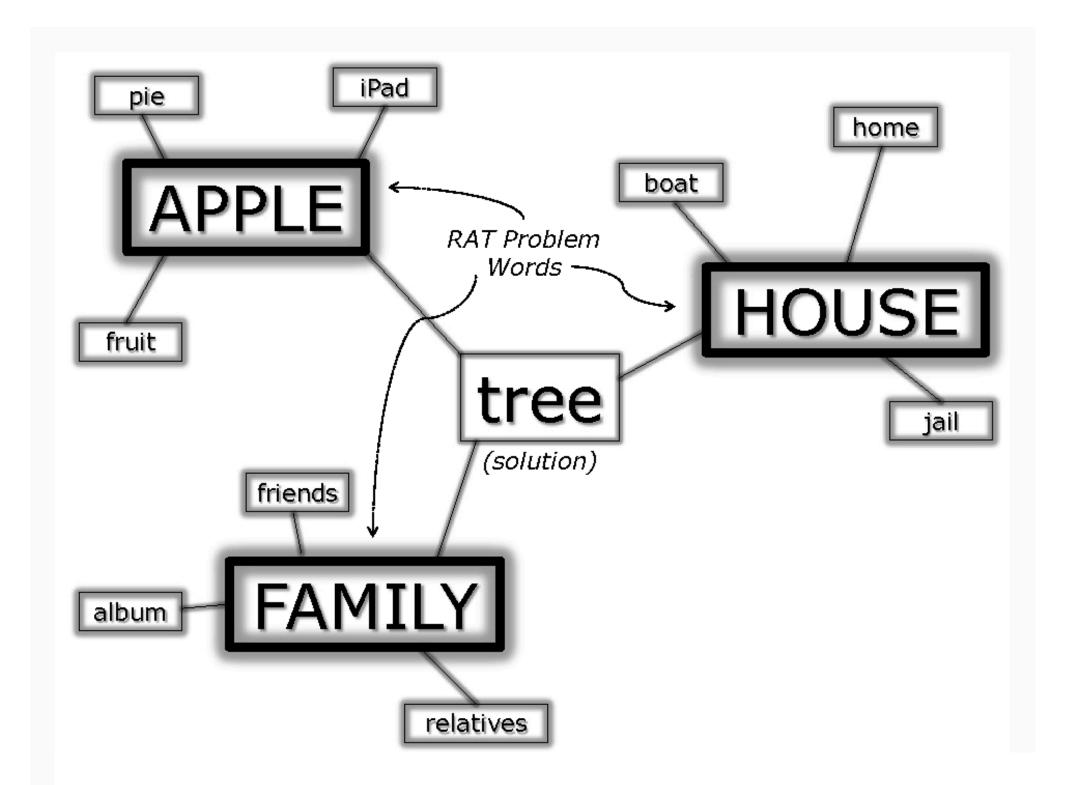
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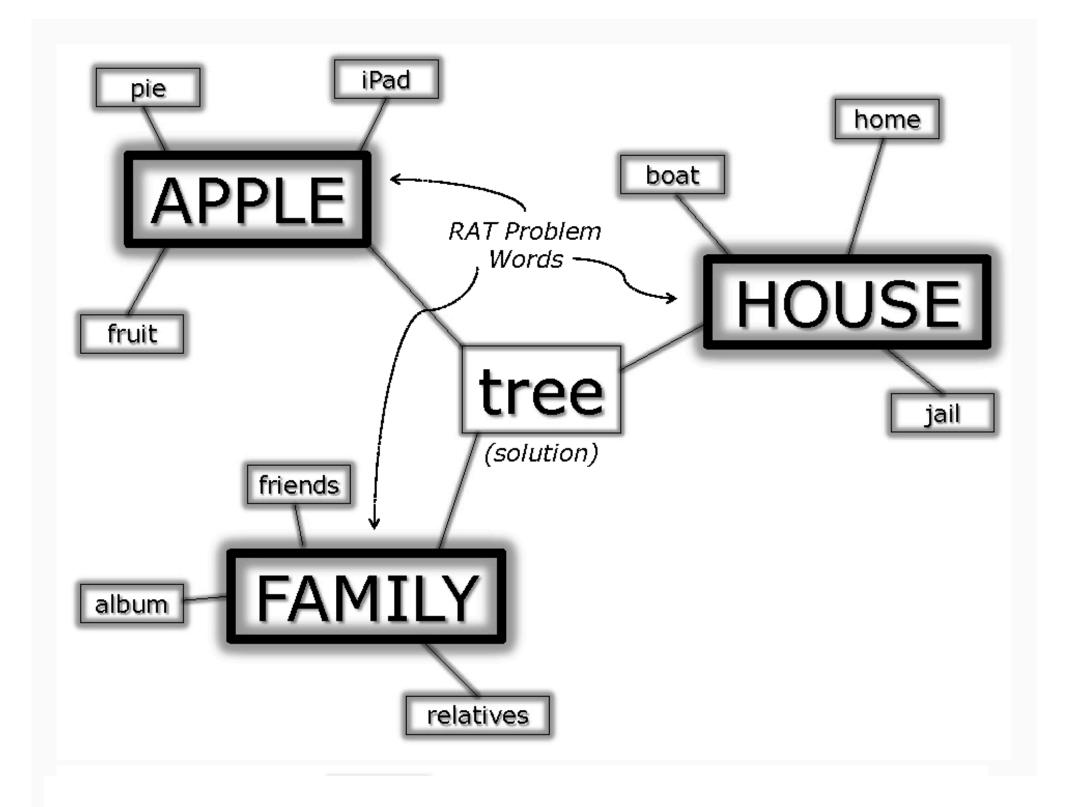
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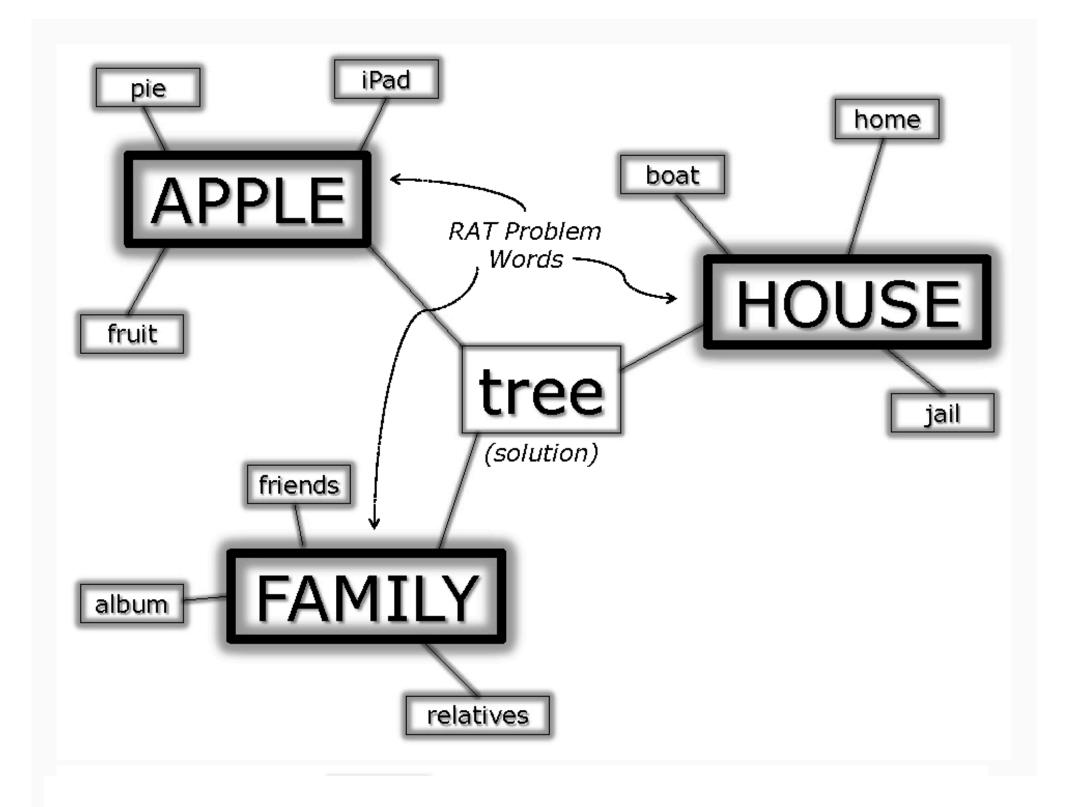
RAT: how does the brain solves it?



Candidates: common nearest neighbor to all three cue words

Constraints from humans: local search, memory constraints, difficulty

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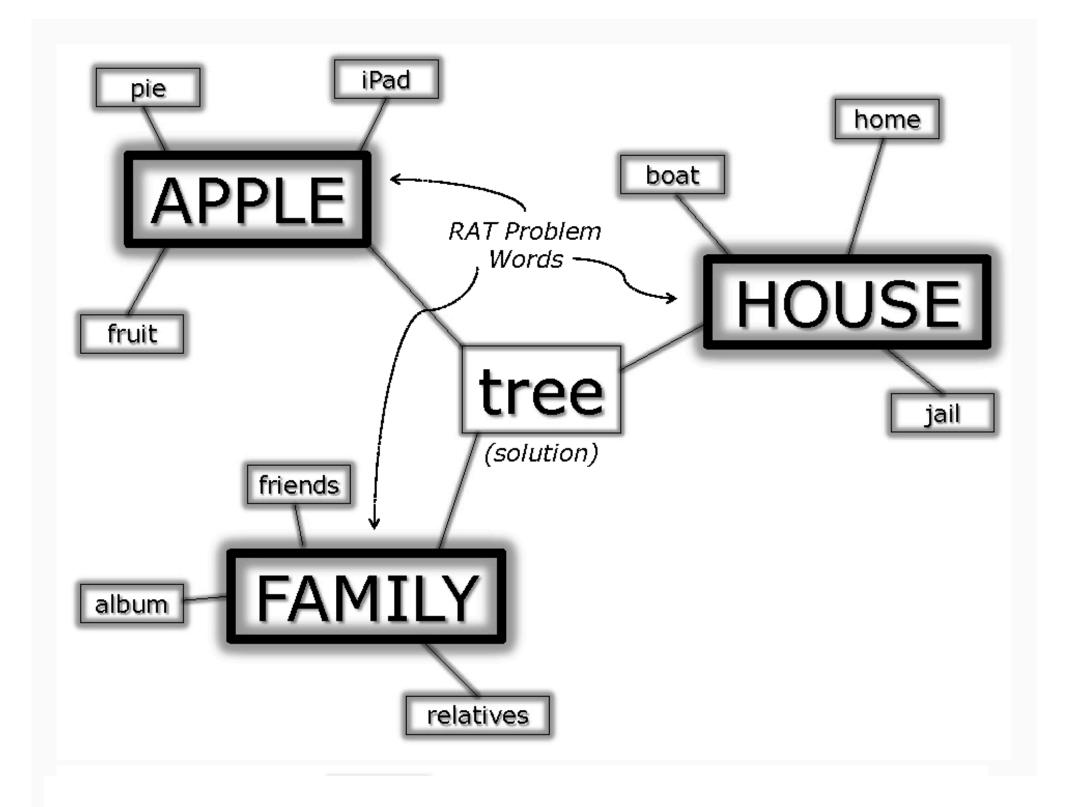


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Naive algorithm 1. Launch random walks from 3 initial cue words Stopping criteria: 1st commonly visited word

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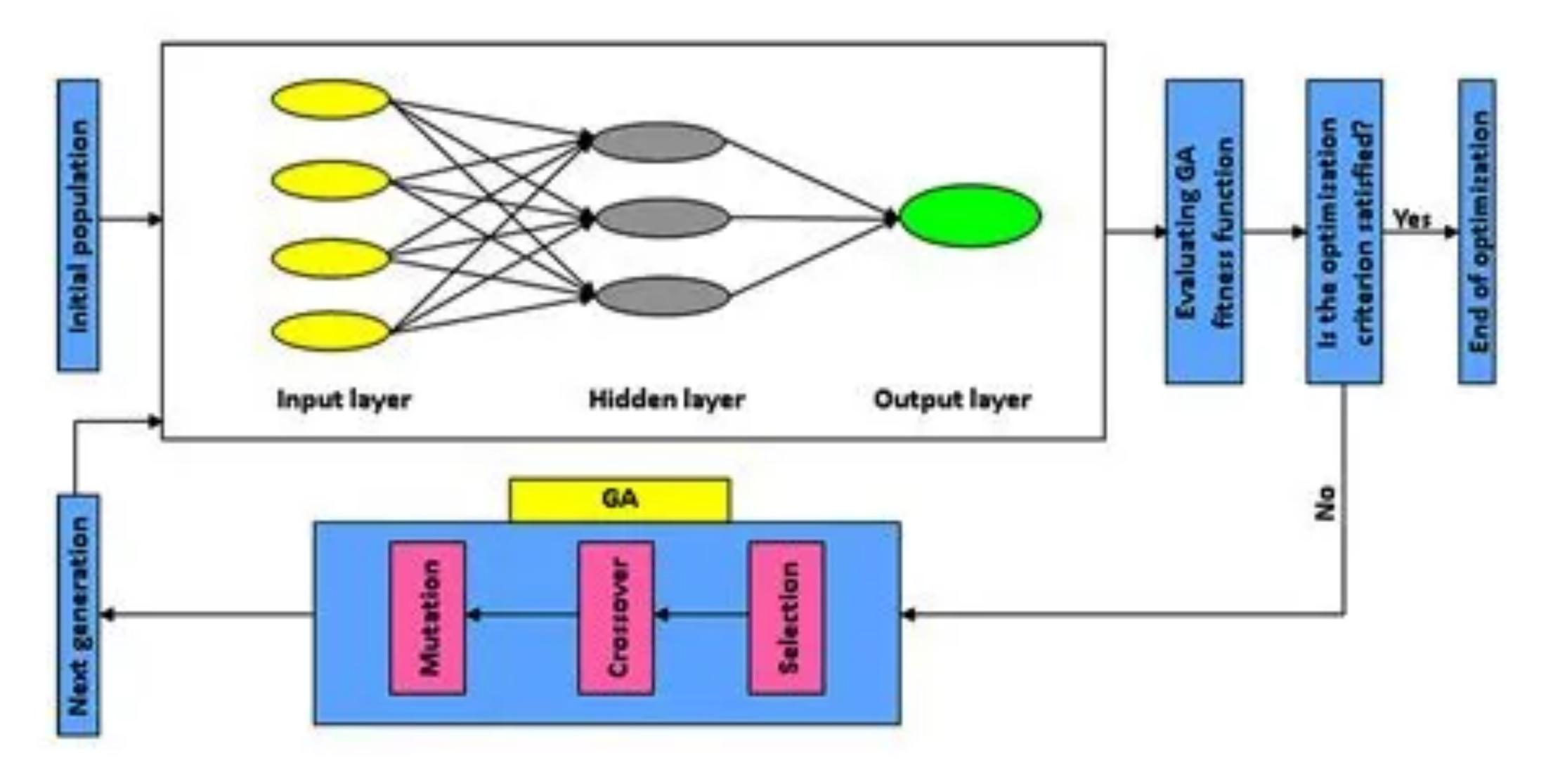
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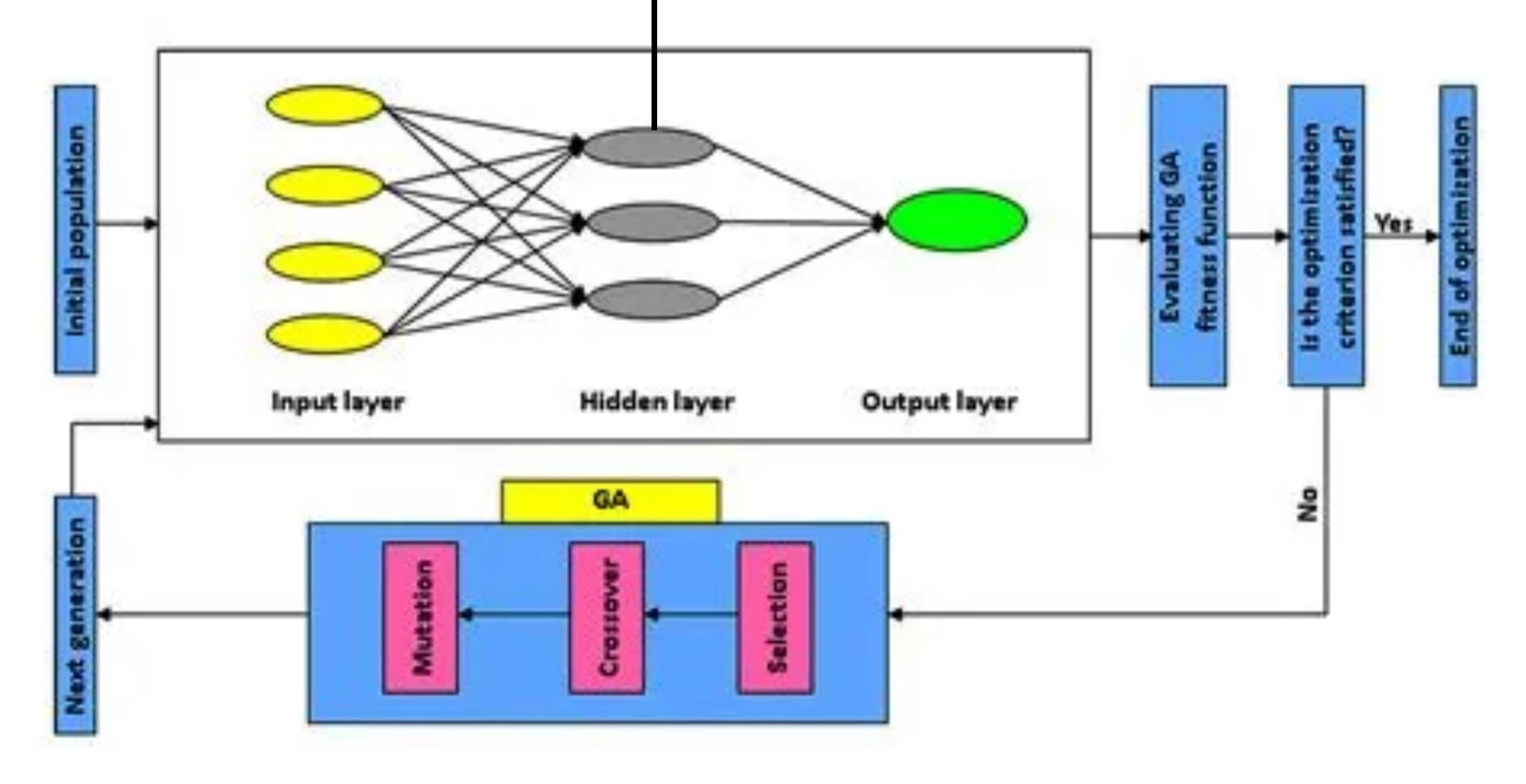
Naive algorithm 2. Represent each word with spectral coordinates Output: Word with MinAvg Euclidean distance to the 3 cues

Functional role of grid cells? More tasks involving grid cells...

GP for grid cells?

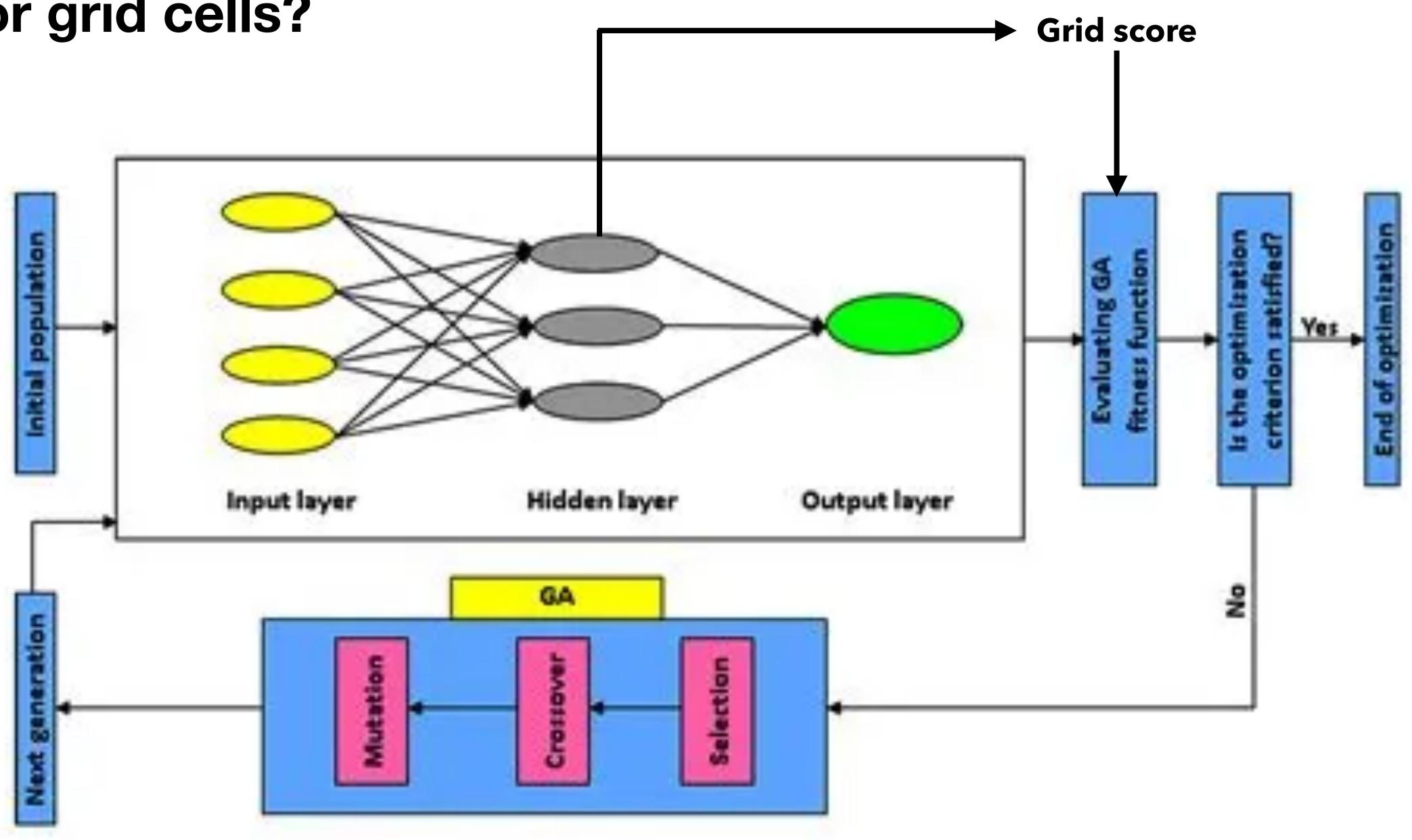


GP for grid cells?



Grid score

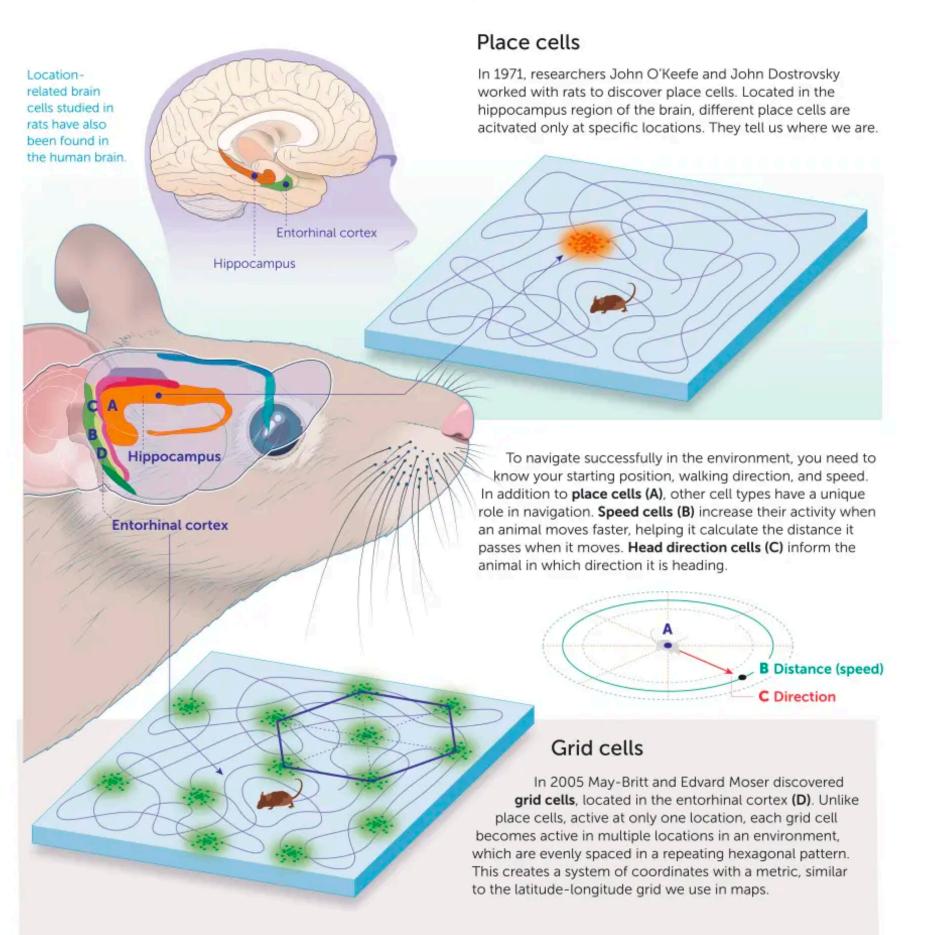
GP for grid cells?



frontiers

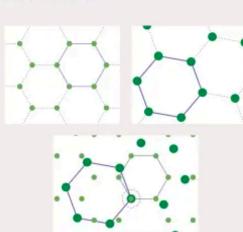
How our body's internal navigation system works

Different types of cells in the brain help us create an internal spatial representation of our environment or "cognitive map" and play a unique role in navigation.



Grid cells help determine location ...

Different grid cells produce different hexagonal patterns at different scales and shifted with respect to the grids of other cells. Using the overlapping grids of several cells, location can be uniquely identified by coincident firing of multiple grid cells.



Location Satellite Satellite

... similarly to GPS systems

In a similar way, in GPS systems the distance measuring of one satelite alone can't pinpoint a location, but the overlap of at last three satelites is used to pinpoint the exact location of a receiver.

position)

- They are thought to allow us to engage **mental navigation** in abstract spaces
- Did the brain discovered Fourier analysis?

Source: How do we Find our Way? Grid Cells in the Brain, by May-Britt Moser. Published by Frontiers for Young Minds, Sep. 2021

Infographic by 5W Infographic

Grid cells periodic pattern may

allow the brain to form cognitive maps in physical space (coordinate system for distance and orientation in space; how things are related to each other independently of the observer

